

# Analytical Modeling of Magnetic Field Distribution in Spoke-Type Permanent-Magnet Machines

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## Abstract:

A two-dimensional (2-D) analytical method for magnetic field and electromagnetic torque calculation in high-speed spoke-type interior permanent-magnet (PM) machines considering slotting effects, magnetization orientation and winding layout has been proposed in this paper. The analytical method is based on the formal resolution of Laplace's and Poisson's equations as well as the Maxwell's equations in polar coordinate by using subdomain technique and applying hyperbolic functions. The proposed method is applied on the performance computation of a prototype spoke-type PM machine (i.e., a 3-phases 18S-8P motor). The analytical results are validated through 2-D finite-element method (FEM) and experimental tests.

**Keywords:** Spoke-type motor, Analytical modeling, Subdomain technique, Numerical, Experimental validation.

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# 1. Introduction

Interior PM machines are interested in industrial applications, especially in the electric vehicle due to their high efficiency, power density and robustness [1]. An accurate prediction of air-gap magnetic field distribution is necessary in order to calculate electromagnetic torque, back electromotive force (EMF) and self- or/and mutual inductances more precise. A variety of techniques including analytical and numerical methods has been conducted to evaluate the magnetic field distribution in electrical machines. Numerical methods like FEM give accurate results and are time consuming specially in first step of design stage [2]-[3]. Analytical methods including conformal mapping [4]-[7], magnetic equivalent circuit [8]-[13], Maxwell/Fourier method [14]-[50] and slot relative permeance calculation [51]-[52] are reported to model electrical machines and are useful in first step of performance evaluation and design optimization stage. The inaccuracy of conformal mapping method in modeling of magnetic field distribution in spoke-type PM machines is due to the presence of a deep and small thickness of PM domain. The subdomain technique is more accurate than the other analytical models [8]. It is interesting to note that the saturation effects can be taken into account of this model type [45]-[46]. This method is based on the formal resolution of Laplace's and Poisson's equations in different regions by applying boundary conditions (BCs) for electrical machines [14]-[50].

To author's knowledge, a few analytical models are presented to calculate magnetic field in spoke-type PM machines using subdomain technique [29]-[33]. The novelty and contribution of the manuscript compared to other works is to propose a straightforward expression for magnetic vector potential by using hyperbolic functions. No references in the literature addressing the issue of experimental verification of analytical model for high-speed spoke-type PM machines were found.

The focus of this paper is to develop an analytical model based on the formal resolution of Laplace's and Poisson's equations in multiphase spoke-type PM machines by using the subdomain technique considering slotting effects, magnetization orientation and winding layout. It is shown that the developed model can effectively estimate magnetic field, electromagnetic torque, back-EMF and self-/mutual inductances. This model is applied on the performance calculation of a prototype spoke-type PM motor (i.e., a 3-phases 18S-8P motor). It is shown that the results of analytical model are in close agreement with the results of FEM and experimental tests.

# 2. Subdomain Definition

The schematic representation of investigated machine is shown in Fig. 1. The machine model is divided into six subdomains. The stator which has two subdomains including  $Q_1$  slot regions (domain  $m$ ) and  $Q_1$  slot opening regions (domain  $l$ ) and the air-gap subdomain (region  $l$ ) are shown in Fig. 2. The rotor has three subdomains including  $Q_2$  inner slot regions (domain  $i$ ),  $Q_2$  PM regions (domain  $j$ ), and  $Q_2$  slot-opening regions (domain  $k$ ), as shown in Fig. 3.

The angular position of the  $i$ -th stator slot,  $i$ -th stator slot-opening,  $i$ -th rotor slot-opening, PM and inner slot are defined as (1), (2), (3) and (4), respectively.

$$\theta_m = -\frac{\alpha}{2} + \frac{2i\pi}{Q_1} \quad \text{with} \quad 1 \leq i \leq Q_1 \quad (1)$$

$$\theta_l = -\frac{\beta}{2} + \frac{2i\pi}{Q_1} \quad \text{with} \quad 1 \leq i \leq Q_1 \quad (2)$$

$$\theta_k = -\frac{\gamma}{2} + \frac{2i\pi}{Q_2} \quad \text{with} \quad 1 \leq i \leq Q_2 \quad (3)$$

$$\theta_i = -\frac{\delta}{2} + \frac{2i\pi}{Q_2} \quad \text{with} \quad 1 \leq i \leq Q_2 \quad (4)$$

The following assumptions are made in theoretical analysis:

- Permeability of rotor and stator cores are infinite;
- End effects are neglected.

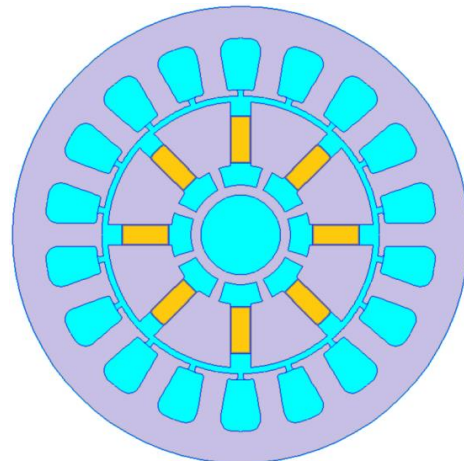


Fig. 1. The geometrical representation of investigated machine.

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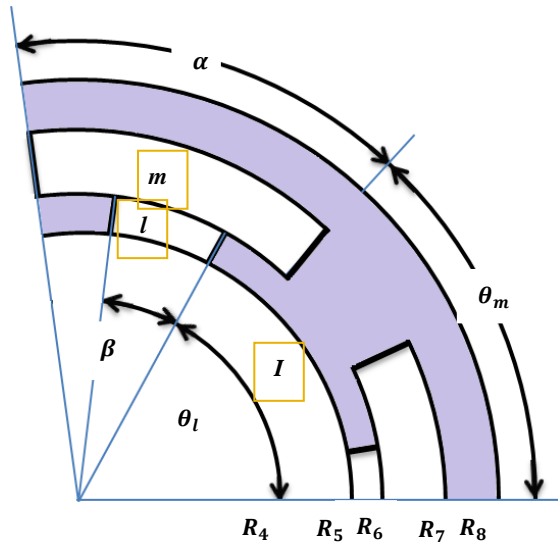


Fig. 2. The stator subdomains including l and m regions.

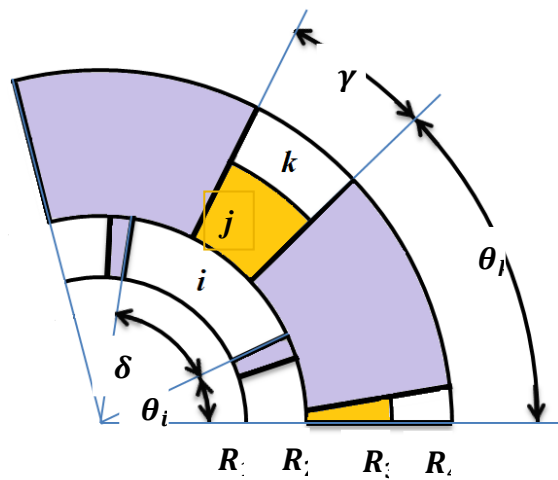


Fig. 3. The rotor subdomains including i, j and k regions.

### 3. Magnetic vector potential calculation

General solution of Laplace's or Poisson's equations in each subdomain is developed in this section. The Laplace equation can be described in polar form as

$$\frac{\partial^2 A}{\partial r^2} + \frac{1}{r} \frac{\partial A}{\partial r} + \frac{1}{r^2} \frac{\partial^2 A}{\partial \theta^2} = 0 \quad \text{for} \quad \begin{cases} R_1 \leq r \leq R_2 \\ \theta_1 \leq \theta \leq \theta_2 \end{cases} \quad (5)$$

Replacing  $r$  by  $R_1 e^{-t}$ , one obtains

$$\frac{\partial^2 A}{\partial t^2} + \frac{\partial^2 A}{\partial \theta^2} = 0 \quad \text{for} \quad \begin{cases} \ln\left(\frac{R_1}{R_2}\right) \leq t \leq 0 \\ \theta_1 \leq \theta \leq \theta_2 \end{cases} \quad (6)$$

### 3.1. Magnetic Vector Potential in the Stator Slot Subdomain (Region m)

The Poisson's equation in the stator inner slot subdomain is given by

$$\frac{\partial^2 A_m}{\partial t^2} + \frac{\partial^2 A_m}{\partial \theta^2} = -\mu_0 J \quad \text{for} \quad \begin{cases} t_1 \leq t \leq t_2 \\ \theta_m \leq \theta \leq \theta_m + \alpha \end{cases} \quad (7)$$

where  $t_1 = \ln\left(\frac{R_6}{R_7}\right)$  and  $t_2 = 0$ .

Neumann BCs at the bottom and at each side of the slot are obtained as

$$\left. \frac{\partial A_i}{\partial \theta} \right|_{\theta=\theta_m} = 0 \quad \text{and} \quad \left. \frac{\partial A_i}{\partial \theta} \right|_{\theta=\theta_m+\alpha} = 0 \quad (8)$$

$$\left. \frac{\partial A_i}{\partial t} \right|_{t=t_1} = 0 \quad (9)$$

The general solution of (7) using the separation of variables method is given by

$$A_m(t, \theta) = a_0^m - \frac{1}{2} \mu_0 J_i \left( e^{-t_1 t} + \frac{1}{2} e^{-2t+t_1} \right) + \sum_{h=1}^{\infty} \left( a_h^m \frac{\alpha}{h\pi} \frac{\text{Cosh}\left(\frac{h\pi}{\alpha}(t-t_1)\right)}{\text{Sinh}\left(\frac{h\pi}{\alpha}(t_2-t_1)\right)} \cdot \text{Cos}\left(\frac{h\pi}{\alpha}(\theta-\theta_m)\right) \right) \quad (10)$$

where  $h$  is a positive integer, the coefficients  $a_0^m$  and  $a_h^m$  are determined based on the continuity and interface conditions and

$$J_i(\theta) = J_{i,0} + \sum_{h=1}^{\infty} J_{i,h} \cos\left(\frac{h\pi}{\alpha}(\theta-\theta_m)\right) \quad (11)$$

with

$$J_{i,0} = \frac{1}{2} (J_{1i} + J_{2i}) \quad (12)$$

$$J_{i,h} = \frac{2}{h\mu} [J_{1i} + (-1)^h J_{2i}] \sin\left(\frac{h\pi}{2}\right) \quad (13)$$

The continuity of the magnetic vector potential between the sub-domain  $m$  and the region  $l$  leads to

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$$\frac{\partial A_m}{\partial t} \Big|_{t=t_2} = f(\theta)$$

$$= \begin{cases} \frac{\partial A_l}{\partial t} \Big|_{t=t_3} & \text{for } \theta_l \leq \theta \leq \theta_l + \beta \\ 0 & \text{elsewhere} \end{cases} \quad (14)$$

Interface condition (14) gives

$$\mu_0 J_i \text{Sinh}(t_1) = \frac{1}{\alpha} \int_{\theta_m}^{\theta_m + \alpha} f(\theta) \cdot d\theta \quad (15)$$

$$a_h^m = \frac{2}{\alpha} \int_{\theta_m}^{\theta_m + \alpha} f(\theta) \cdot \text{Cos} \left( \frac{h\pi}{\alpha} (\theta - \theta_m) \right) \cdot d\theta \quad (16)$$

### 3.2. Magnetic Vector Potential in the Stator Slot-Opening Subdomain (Region I)

The Laplace's equation in the stator second inner slot-opening subdomain is given by

$$\frac{\partial^2 A_l}{\partial t^2} + \frac{\partial^2 A_l}{\partial \theta^2} = 0 \quad \text{for } \begin{cases} t_3 \leq t \leq t_4 \\ \theta_l \leq \theta \leq \theta_l + \beta \end{cases} \quad (17)$$

where  $\mathbf{t}_3 = \ln \left( \frac{R_5}{R_6} \right)$  and  $\mathbf{t}_4 = \mathbf{0}$ .

Neumann BCs at the bottom and at each side of the slot are obtained as

$$\frac{\partial A_l}{\partial \theta} \Big|_{\theta=\theta_l} = 0 \quad \text{and} \quad \frac{\partial A_l}{\partial \theta} \Big|_{\theta=\theta_l+\beta} = 0 \quad (18)$$

The general solution of (17) using the separation of variables method is given by

$$A_l(t, \theta) = a_0^1 + b_0^1 t$$

$$+ \sum_{h=1}^{\infty} \left( \frac{\text{Sinh} \left( \frac{h\pi}{\beta} (t - t_4) \right)}{\text{Sinh} \left( \frac{h\pi}{\beta} (t_3 - t_4) \right)} a_h^1 + \frac{\text{Sinh} \left( \frac{h\pi}{\beta} (t - t_3) \right)}{\text{Sinh} \left( \frac{h\pi}{\beta} (t_4 - t_3) \right)} b_h^1 \right) \cdot \text{Cos} \left( \frac{h\pi}{\beta} (\theta - \theta_l) \right) \quad (19)$$

where  $m$  is a positive integer and the coefficients,  $a_0^1$ ,  $b_0^1$ ,  $a_h^1$  and  $b_h^1$  are determined based on the continuity and interface conditions.

The continuity of the magnetic vector potential between the subdomain I and the regions  $m$  and I leads to

$$A_l(t_4, \theta) = A_l(t_5, \theta) \quad \text{for } \theta_l \leq \theta \leq \theta_l + \beta \quad (20)$$

$$A_l(t_3, \theta) = A_m(t_4, \theta) \quad \text{for } \theta_l \leq \theta \leq \theta_l + \beta \quad (21)$$

Interface condition (20) gives

$$a_0^1 = \frac{1}{\beta} \int_{\theta_l}^{\theta_l + \beta} A_l(t_5, \theta) \cdot d\theta \quad (22)$$

$$b_m^1 = \frac{2}{\beta} \int_{\theta_l}^{\theta_l + \beta} A_l(t_5, \theta) \cdot \text{Cos} \left( \frac{h\pi}{\beta} (\theta - \theta_l) \right) \cdot d\theta \quad (23)$$

Interface condition (21) gives

$$a_0^1 + \ln \left( \frac{R_5}{R_6} \right) b_0^1 = \frac{1}{\beta} \int_{\theta_l}^{\theta_l + \alpha} A_m(t_4, \theta) \cdot d\theta \quad (24)$$

$$a_h^1 = \frac{2}{\beta} \int_{\theta_l}^{\theta_l + \beta} A_m(t_4, \theta) \cdot \text{Cos} \left( \frac{h\pi}{\beta} (\theta - \theta_l) \right) \cdot d\theta \quad (25)$$

### 3.3. Magnetic Vector Potential in the Air-gap Subdomain (Region I)

The Laplace's equation in the internal air-gap subdomain is given by

$$\frac{\partial^2 A_l}{\partial t^2} + \frac{\partial^2 A_l}{\partial \theta^2} = 0 \quad \text{for } \begin{cases} t_5 \leq t \leq t_6 \\ 0 \leq \theta \leq 2\pi \end{cases} \quad (26)$$

where  $\mathbf{t}_5 = \ln \left( \frac{R_4}{R_5} \right)$  and  $\mathbf{t}_6 = \mathbf{0}$ .

The general solution of (26) considering periodicity boundary conditions is obtained as

$$A_l(t, \theta) = \sum_{n=1}^{\infty} \left( \frac{1 \cdot \text{Cosh}(n(t - t_6))}{n \cdot \text{Sinh}(n(t_5 - t_6))} a_n^1 + \frac{1 \cdot \text{Cosh}(n(t - t_5))}{n \cdot \text{Sinh}(n(t_6 - t_5))} b_n^1 \right) \text{Cos}(n\theta)$$

$$+ \sum_{n=1}^{\infty} \left( \frac{1 \cdot \text{Cosh}(n(t - t_6))}{n \cdot \text{Sinh}(n(t_5 - t_6))} c_n^1 + \frac{1 \cdot \text{Cosh}(n(t - t_5))}{n \cdot \text{Sinh}(n(t_6 - t_5))} d_n^1 \right) \text{Sin}(n\theta) \quad (27)$$

where  $n$  is a positive integer.

The coefficients  $a_n^I$ ,  $b_n^I$ ,  $c_n^I$  and  $d_n^I$  are determined considering the continuity of magnetic vector potential between the internal air-gap subdomain I and the region I using a Fourier series expansion of interface condition (28) and (29) over the air-gap interval.

The continuity of the magnetic vector potential between the internal air-gap subdomain I and the regions l and k leads to

$$\frac{\partial A_I}{\partial t} \Big|_{t=t_5} = g(\theta) = \begin{cases} \frac{\partial A_I}{\partial t} \Big|_{t=t_4} & \text{for } \theta_l \leq \theta \leq \theta_l + \beta \\ 0 & \text{elsewhere} \end{cases} \quad (28)$$

$$\frac{\partial A_I}{\partial t} \Big|_{t=t_6} = h(\theta) = \begin{cases} \frac{\partial A_k}{\partial t} \Big|_{t=t_7} & \text{for } \theta_k \leq \theta \leq \theta_k + \gamma \\ 0 & \text{elsewhere} \end{cases} \quad (29)$$

Interface condition (28) gives

$$a_n^I = \frac{2}{2\pi} \int_{\theta_l}^{\theta_l + \beta} g(\theta) \cdot \text{Cos}(n\theta) \cdot d\theta \quad (27)$$

$$c_n^I = \frac{2}{2\pi} \int_{\theta_l}^{\theta_l + \beta} g(\theta) \cdot \text{Sin}(n\theta) \cdot d\theta \quad (28)$$

Interface condition (29) gives

$$b_n^I = \frac{2}{2\pi} \int_{\theta_k}^{\theta_k + \gamma} h(\theta) \cdot \text{Cos}(n\theta) \cdot d\theta \quad (29)$$

$$d_n^I = \frac{2}{2\pi} \int_{\theta_k}^{\theta_k + \gamma} h(\theta) \cdot \text{Sin}(n\theta) \cdot d\theta \quad (30)$$

### 3.4. Magnetic Vector Potential in the Rotor Slot-Opening Subdomain (Region k)

The Laplace's equation in the stator second inner slot-opening subdomain is given by

$$\frac{\partial^2 A_k}{\partial t^2} + \frac{\partial^2 A_k}{\partial \theta^2} = 0 \quad \text{for } \begin{cases} t_7 \leq t \leq t_8 \\ \theta_k \leq \theta \leq \theta_k + \gamma \end{cases} \quad (34)$$

where  $t_7 = \ln\left(\frac{R_3}{R_4}\right)$  and  $t_8 = 0$ .

Neumann BCs at the bottom and at each side of the slot are obtained as

$$\frac{\partial A_{II}}{\partial \theta} \Big|_{\theta=\theta_k} = 0 \quad \text{and} \quad \frac{\partial A_{II}}{\partial \theta} \Big|_{\theta=\theta_k + \gamma} = 0 \quad (35)$$

The general solution of (34) using the separation of variables method is given by

$$A_k(t, \theta) = a_0^k + b_0^k t + \sum_{h=1}^{\infty} \left( \frac{\text{Sinh}\left(\frac{h\pi}{\gamma}(t - t_8)\right)}{\text{Sinh}\left(\frac{h\pi}{\gamma}(t_7 - t_8)\right)} a_h^k + \frac{\text{Sinh}\left(\frac{h\pi}{\gamma}(t - t_7)\right)}{\text{Sinh}\left(\frac{h\pi}{\gamma}(t_8 - t_7)\right)} b_h^k \right) \cdot \text{Cos}\left(\frac{h\pi}{\gamma}(\theta - \theta_k)\right) \quad (36)$$

where h is a positive integer and the coefficients,  $a_0^k$ ,  $b_0^k$ ,  $a_h^k$  and  $b_h^k$  are determined based on the continuity and interface conditions.

The continuity of the magnetic vector potential between the subdomain k and the regions j and I leads to

$$A_k(t_8, \theta) = A_j(t_9, \theta) \quad \text{for } \theta_k \leq \theta \leq \theta_k + \gamma \quad (37)$$

$$A_k(t_7, \theta) = A_I(t_6, \theta) \quad \text{for } \theta_k \leq \theta \leq \theta_k + \gamma \quad (38)$$

Interface condition (37) gives

$$a_0^k = \frac{1}{\gamma} \int_{\theta_k}^{\theta_k + \gamma} A_j(t_9, \theta) \cdot d\theta \quad (39)$$

$$b_h^k = \frac{2}{\gamma} \int_{\theta_k}^{\theta_k + \gamma} A_j(t_9, \theta) \cdot \text{Cos}\left(\frac{h\pi}{\gamma}(\theta - \theta_k)\right) \cdot d\theta \quad (40)$$

Interface condition (38) gives

$$a_0^k + \ln\left(\frac{R_3}{R_4}\right) b_0^k = \frac{1}{\gamma} \int_{\theta_l}^{\theta_l + \gamma} A_I(t_6, \theta) \cdot d\theta \quad (41)$$

$$a_h^k = \frac{2}{\gamma} \int_{\theta_l}^{\theta_l + \gamma} A_I(t_6, \theta) \cdot \text{Cos}\left(\frac{h\pi}{\gamma}(\theta - \theta_k)\right) \cdot d\theta \quad (42)$$

### 3.5. Magnetic Vector Potential in the PM Subdomain (Region j)

The Poisson's equation in the stator PM subdomain is given by

$$\frac{\partial^2 A_j}{\partial t^2} + \frac{\partial^2 A_j}{\partial \theta^2} = -\mu_0 \frac{M_\theta}{r} \quad (43)$$

for  $\begin{cases} t_9 \leq t \leq t_{10} \\ \theta_k \leq \theta \leq \theta_k + \gamma \end{cases}$

where  $t_9 = \ln\left(\frac{R_2}{R_3}\right)$  and  $t_{10} = 0$ .

Neumann BCs at the bottom and at each side of the slot are obtained as

$$\left. \frac{\partial A_j}{\partial \theta} \right|_{\theta=\theta_k} = 0 \quad \text{and} \quad \left. \frac{\partial A_j}{\partial \theta} \right|_{\theta=\theta_k+\gamma} = 0 \quad (44)$$

The general solution of (40) is written as

$$A_j(t, \theta) = a_0^j + b_0^j t - \mu_0 M_j R_2 e^{-t} + \sum_{h=1}^{\infty} \left( \frac{\gamma \cosh\left(\frac{h\pi}{\gamma}(t-t_{10})\right)}{h\pi \sinh\left(\frac{h\pi}{\gamma}(t_9-t_{10})\right)} a_h^j + \frac{\gamma \cosh\left(\frac{h\pi}{\gamma}(t-t_9)\right)}{h\pi \sinh\left(\frac{h\pi}{\gamma}(t_{10}-t_9)\right)} b_h^j \right) \cos\left(\frac{h\pi}{\gamma}(\theta-\theta_k)\right) \quad (45)$$

where  $M_\theta = M_j = (-1)^i B_r / \mu_0$ ,  $l$  is a positive integer and the coefficients  $a_0^j$ ,  $b_0^j$ ,  $a_h^j$  and  $b_h^j$  are determined based on the continuity and interface conditions.

The continuity of the magnetic vector potential between the subdomain j and the region k and i leads to

$$\left. \frac{\partial A_j}{\partial t} \right|_{t=t_9} = \left. \frac{\partial A_k}{\partial t} \right|_{t=t_8} \quad \text{for} \quad \theta_k \leq \theta \leq \theta_k + \gamma \quad (46)$$

$$\left. \frac{\partial A_j}{\partial t} \right|_{t=t_{10}} = \left. \frac{\partial A_i}{\partial t} \right|_{t=t_{11}} \quad \text{for} \quad \theta_k \leq \theta \leq \theta_k + \gamma \quad (47)$$

Interface condition (46) gives

$$b_0^j + \mu_0 M_j R_2 e^{-t_9} = \frac{1}{\gamma} \int_{\theta_k}^{\theta_k+\gamma} \frac{\partial A_k}{\partial t} \Big|_{t=t_8} \cdot d\theta \quad (48)$$

$$a_h^j = \frac{2}{\gamma} \int_{\theta_k}^{\theta_k+\gamma} \frac{\partial A_k}{\partial t} \Big|_{t=t_8} \cdot \cos\left(\frac{h\pi}{\gamma}(\theta-\theta_k)\right) \cdot d\theta \quad (49)$$

$$\cos\left(\frac{h\pi}{\gamma}(\theta-\theta_k)\right) \cdot d\theta$$

Interface condition (44) gives

$$b_0^j + \mu_0 M_j R_2 e^{-t_{10}} = \frac{1}{\gamma} \int_{\theta_k}^{\theta_k+\gamma} \frac{\partial A_i}{\partial t} \Big|_{t=t_{11}} \cdot d\theta \quad (50)$$

$$b_h^j = \frac{2}{\gamma} \int_{\theta_k}^{\theta_k+\gamma} \frac{\partial A_i}{\partial t} \Big|_{t=t_{11}} \cdot \cos\left(\frac{h\pi}{\gamma}(\theta-\theta_k)\right) \cdot d\theta \quad (51)$$

### 3.6. Magnetic Vector Potential in the Subdomain (Region i)

The Laplace equation in the i-th rotor slot subdomain is given by

$$\frac{\partial^2 A_i}{\partial t^2} + \frac{\partial^2 A_i}{\partial \theta^2} = 0 \quad \text{for} \quad \begin{cases} t_{11} \leq t \leq t_{12} \\ \theta_i \leq \theta \leq \theta_i + \delta \end{cases} \quad (52)$$

The general solution of (52) using the separation of variables method based on boundary conditions is

$$A_i(t, \theta) = \sum_{h=1}^{\infty} \left( a_h^i \frac{\cosh\left(\frac{h\pi}{\gamma}(t-t_{12})\right)}{\cosh\left(\frac{h\pi}{\gamma}(t_{11}-t_{12})\right)} \right) \cos\left(\frac{h\pi}{\delta}(\theta-\theta_i)\right) \quad (53)$$

where  $k$  is a positive integer and  $a_h^i$  is an arbitrary constants.

The continuity of the magnetic vector potential between the i-th slot and the internal air-gap regions I leads to

$$A_i(t_{11}, \theta) = A_j(t_{10}, \theta) \quad \text{for} \quad \theta_i \leq \theta \leq \theta_i + \delta \quad (54)$$

Considering interface condition (54), the integration constants can be determined as

$$a_h^i = \frac{2}{\gamma} \int_{\theta_i}^{\theta_i+\gamma} A_j(t_{10}, \theta) \cdot \cos\left(\frac{h\pi}{\delta}(\theta-\theta_i)\right) \cdot d\theta \quad (55)$$



## 4. Performance Calculation and Model Evaluation

### 4.1. Performance Computation

The electromagnetic torque and cogging torque components are obtained using the Maxwell stress tensor in and expressed as

$$T_e = \frac{L_s}{\mu_0} \int_0^{2\pi} B_{1r}(t_e, \theta) \cdot B_{1\theta}(t_e, \theta) \cdot d\theta \quad (56)$$

$$B_{1r}(t_e, \theta) = -\frac{e^{t_e}}{R_2} \left( \sum_{n=1}^{\infty} \left( \frac{a_n^l \cosh(n(t_e - t_6))}{\sinh(n(t_5 - t_6))} + b_n^l \frac{\cosh(n(t_e - t_5))}{\sinh(n(t_6 - t_5))} \right) \cdot \sin(n\theta) - \sum_{n=1}^{\infty} \left( \frac{c_n^l \cosh(n(t_e - t_6))}{\sinh(n(t_5 - t_6))} + d_n^l \frac{\cosh(n(t_e - t_5))}{\sinh(n(t_6 - t_5))} \right) \cdot \cos(n\theta) \right) \quad (57)$$

$$B_{1\theta}(t_e, \theta) = -\frac{e^{t_e}}{R_2} \left( \sum_{n=1}^{\infty} \left( \frac{a_n^l \sinh(n(t_e - t_6))}{\sinh(n(t_5 - t_6))} + b_n^l \frac{\sinh(n(t_e - t_5))}{\sinh(n(t_6 - t_5))} \right) \cdot \cos(n\theta) + \sum_{n=1}^{\infty} \left( \frac{c_n^l \sinh(n(t_e - t_6))}{\sinh(n(t_5 - t_6))} + d_n^l \frac{\sinh(n(t_e - t_5))}{\sinh(n(t_6 - t_5))} \right) \cdot \sin(n\theta) \right) \quad (58)$$

where  $L_s$  is the axial length of the motor and  $t_e$  is calculated by

$$t_e = \ln\left(\frac{R_4}{R_e}\right) \text{ with } R_e = (R_4 + R_5)/2 \quad (59)$$

For double layer winding, the phase flux vector is calculated by

$$\begin{bmatrix} \psi_a \\ \psi_b \\ \psi_c \end{bmatrix} = \begin{bmatrix} \psi_{1a} \\ \psi_{1b} \\ \psi_{1c} \end{bmatrix} + \begin{bmatrix} \psi_{2a} \\ \psi_{2b} \\ \psi_{2c} \end{bmatrix} \quad (60)$$

where

$$\begin{bmatrix} \psi_{1a} \\ \psi_{1b} \\ \psi_{1c} \end{bmatrix} = \frac{N_c}{2} C_1^T [\varphi_{11} \quad \varphi_{12} \quad \varphi_{13} \quad \dots \quad \varphi_{1Q_2}] \quad (61)$$

and

$$\begin{bmatrix} \psi_{2a} \\ \psi_{2b} \\ \psi_{2c} \end{bmatrix} = \frac{N_c}{2} C_2^T [\varphi_{21} \quad \varphi_{22} \quad \varphi_{23} \quad \dots \quad \varphi_{2Q_2}] \quad (62)$$

For the stator slots,  $\varphi$  is given by

$$\varphi_{1i} = -\frac{2L_s R_4^2}{k_f S} \int_0^{\frac{\alpha}{2}} \int_0^{t_5} A_{m_i}(t, \theta) \cdot e^{-2t} \cdot dt \cdot d\theta \quad (63)$$

$$\varphi_{2i} = -\frac{2L_s R_4^2}{k_f S} \int_{\frac{\alpha}{2}}^{\alpha} \int_0^{t_5} A_{m_i}(t, \theta) \cdot e^{-2t} \cdot dt \cdot d\theta \quad (64)$$

The back-EMF of phase  $a$  is given by

$$E_a = \omega \frac{d\psi_a}{d\theta_r} \quad (65)$$

where  $\omega$  is the rotor angular speed and  $\psi_a$  is flux linkage per phase  $a$ .

The stator inductances (self-inductances) of phase  $a$  is given by

$$L = \frac{\psi_a}{I_a} \quad (66)$$

where  $I_a$  is the peak current in phase  $a$ .

### 4.2. Model Evaluation

The investigated motor parameters are given in Table 1. A schematic diagram of the double layer winding of the motor is shown in Fig. 4. The 2-D FEM is applied on performance calculation of the motor. Magnetic field distribution in the motor is represented in Fig. 5. The fabricated spoke-type PM motor and experimental test setup are shown in Fig. 6 and Fig. 7, respectively.

**Table 1 Parameters of the studied motor.**

Symbol	Quantity Unit: angles (Mech. Degree) Dimensions (mm)	Value
$R_1$	Inner radius of the rotor inner slot	12.55
$R_2$	Inner radius of the rotor permanent magnet	18
$R_3$	Inner radius of the rotor slot-opening	30
$R_4$	Inner radius of the airgap	33.5
$R_5$	Inner radius of the stator slot-opening	35
$R_6$	Inner radius of the stator slot	36.5
$R_7$	Outer radius of the stator slot	49.5
$R_8$	Outer radius of the stator yoke	57.5
$\theta_i$	Angular position of the first rotor inner slot	27
$\theta_k$	Angular position of the first rotor slot-opening	22
$\theta_l$	Angular position of the first stator slot-opening	18.77

$\theta_m$	Angular position of the first stator slot	8.88
$\alpha$	The stator slot angle	7.2
$\beta$	The stator slot-opening angle	2.45
$\gamma$	The rotor slot-opening angle	8.55
$\delta$	The rotor inner slot angle	36
$p$	Pole pairs-number	4
$Q_s$	Number of stator slots	18
$B_r$	Remanence of the PM (T)	1.2
$L_s$	Axial length	48

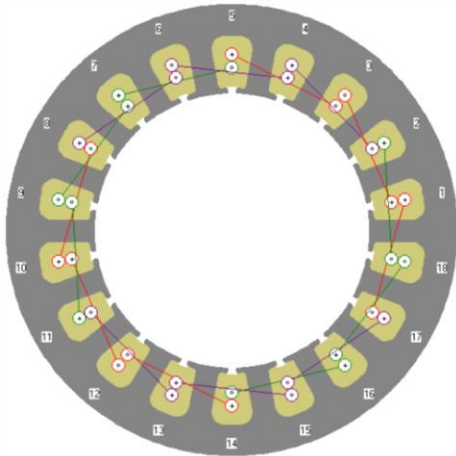


Fig. 4. A schematic diagram of the double layer winding of the investigated motor.

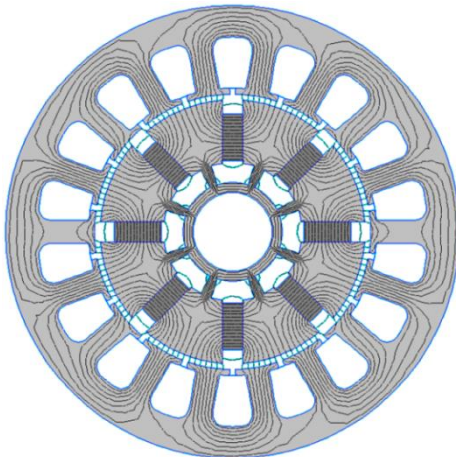


Fig. 5. Magnetic field distribution in 3-phases spoke-type PM motor.

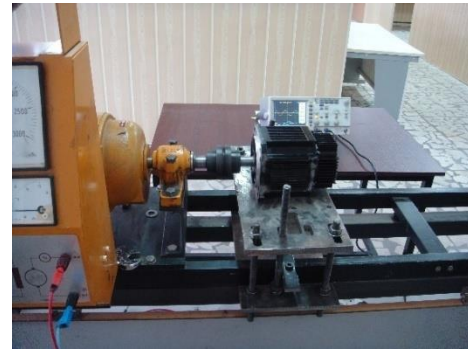


Fig. 6. Experimental test setup.

The motor performance results are compared analytically, numerically and experimentally in open circuit and full load conditions. The open circuit comparison of radial magnetic flux density and cogging torque waveforms are shown in Fig. 8 and Fig. 9, respectively.

The full load results at 6,000 rpm, including electromagnetic torque, back-EMF, self- and mutual inductances waveforms are shown in Fig. 8, Fig. 9, Fig. 10, Fig. 11, Fig. 12, and Fig. 13, respectively.

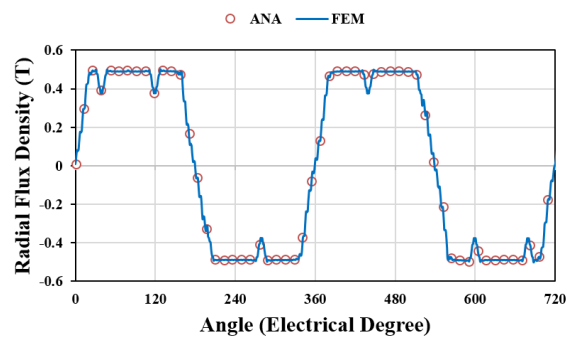


Fig. 7. Analytical and numerical comparison of radial magnetic flux density waveforms in mean radius of the studied motor.

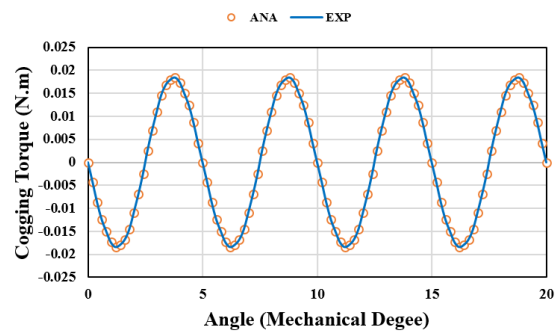


Fig. 8. Analytical and experimental comparison of cogging torque waveforms.



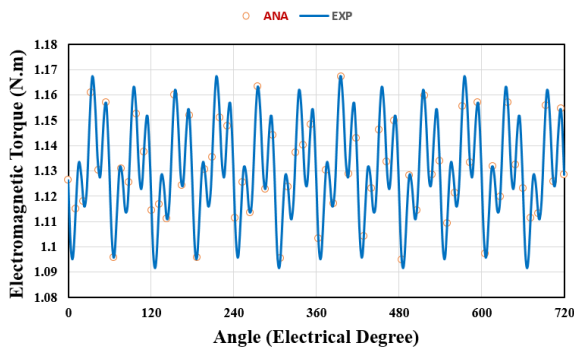


Fig. 9. Analytical and experimental comparison of electromagnetic torque waveforms at 6,000 rpm.

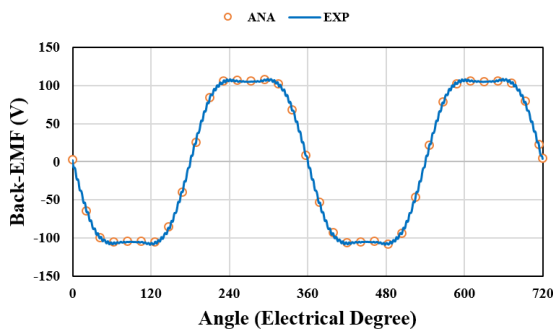


Fig. 10. Analytical and experimental comparison of phase back-EMF waveforms at 6,000 rpm.

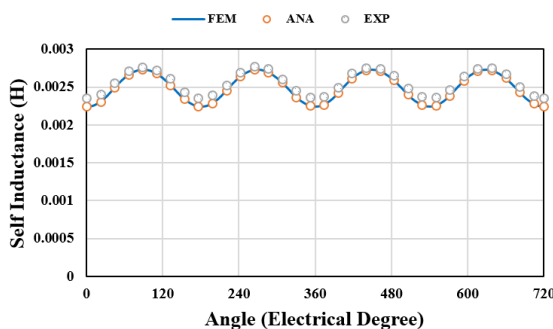


Fig. 11. Analytical and numerical comparison of phase self-inductance waveforms at 6,000 rpm.

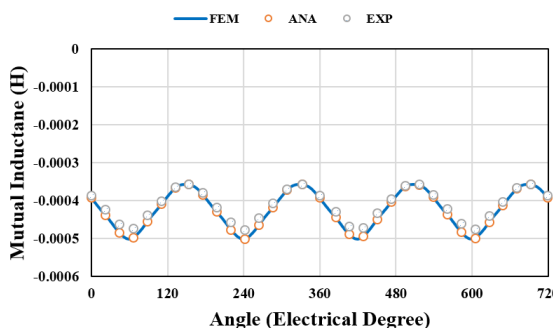


Fig. 12. Analytical and numerical comparison of mutual inductances waveforms at 6,000 rpm.

## 5. Conclusion

A 2-D analytical model for performance prediction in multiphase high-speed spoke-type PM machines considering slotting effects, magnetization orientation and winding layout has been developed in this paper. Fourier analysis method based on the subdomain technique is applied to derive analytical expressions for calculation of magnetic vector potential, magnetic flux density, electromagnetic torque, back-EMF and self- and mutual inductances in spoke-type PM machines. This model is applied for performance computation of a prototype spoke-type PM motor and the results of proposed model are verified thanks to FEM and experimental results.

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