

# 1 Cost-efficient and high precision method for 2 the assembly of LN-based photonic crystal 3 slabs on the fiber tip for the implementation of 4 E-field sensors

5 BRUNO ROBERT, VENANCIO CALERO, MIGUEL-ANGEL SUAREZ, ROLAND  
6 SALUT, FLORENT BEHAGUE, FADI BAIDA, NADEGE COURJAL, MARIA-PILAR  
7 BERNAL\*

8 *FEMTO-ST Institute, University of Bourgogne Franche-Comte, CNRS UMR 6174, Besançon, (France)*

9 *\*mp.bernal@femto-st.fr*

10 **Abstract:** Lab-on-Fiber technology is an emerging topic for sensing cutting-edge technologies  
11 due to the high versatility and functionality that it offers when it is combined with different  
12 sensitive materials. A particular configuration, which consists of the integration of  
13 nanophotonic structures into the tip of a pigtailed fiber, allows the exploitation of light  
14 localization performances to produce highly performing sensors. However, integrating such  
15 tiny structures into the fiber facet requires complex and expensive procedures. In this work, we  
16 report a novel high precision assembly procedure which ensures the parallelism between the  
17 photonic chip and the fiber surface, in addition to the alignment with the light injection into the  
18 nanostructure. The integrated structure consists of an ultra-compact ( $19\ \mu\text{m} \times 19\ \mu\text{m}$ ) Photonic  
19 Crystal Slab (PCS) structure based on a 700 nm thin film of lithium niobate (LN) which is  
20 sensitive to external E-fields via the electro-optic effect. Thus, the assembled sensor detects  
21 electric fields, presenting a great linearity and a sensitivity of 170 V/m. This technique shows  
22 a way to assemble compact planar nanostructures into fiber facets keeping high throughput,  
23 high precision and relatively low costs.

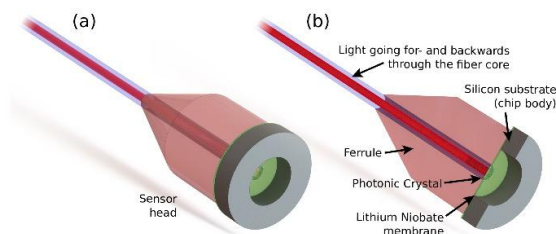
## 24 1. Introduction

25 Optical fibers are experiencing an important growth, in addition to data transport, also in the  
26 sensing field due to their unique versatility. They are nowadays employed in applications such  
27 as non-invasive medical surgery and structures damage monitoring [1,2]. Moreover, the  
28 measurement of other magnitudes such as temperature, refractive index, displacement, pressure  
29 and acceleration have been also demonstrated through the so-called Lab-on-Fiber technologies  
30 [3]. Recently, the progress in micro- and nanofabrication techniques is making possible the  
31 growth of fiber tip sensors whose functional area consist on photonic nanostructures integrated  
32 on the fiber extremity such as photonic circuits [4], allowing a new era the of ultra-compact,  
33 versatile sensors. Moreover, Guided Resonances (GRs) have been recently implemented on the  
34 fiber tip [5,6] due to the coupling efficiency of the fiber propagating mode with the photonic  
35 crystal (PhC) mode, in addition to the light localization and vertical confinement performances  
36 allowing light-matter interaction enhancement [7,8]. In a recent work from our group [9], a GR-  
37 based fiber tip configuration is combined with lithium niobate (LN) in order to enhance its  
38 electro- optical performances producing an ultra-compact and distortion-free electric field  
39 sensing probe. Thus, a sub-micrometer free standing slab with a photonic crystal (PhC) is  
40 designed in a 700 nm TFLN membrane so that it can be integrated on the fiber tip. Despite the  
41 fact that the integration of bulk ferroelectric crystals on a fiber tip is a well-known approach  
42 [10] producing all-dielectric E-field sensing probes, the fiber tip approach based on the GRs  
43 has footprints that are thousands of times smaller than these massive crystal devices. The  
44 interest of reduced dimensions is very high since that, in addition to a submillimeter footprint,  
45 it offers clear advantages in terms of enhanced linearity with respect to the E-field, micrometer  
46 spatial resolution as well as extended bandwidth, only limited by the LN intrinsic bandwidth

47 which is in the order of tens of THz. However, the nanofabrication of LN-based PhC with  
48 nanometric precision on LN is still limited to the use of the Focused Ion Beam (FIB) technology  
49 [11] or e-beam assisted IBEE techniques [12], which supposes a constraint on the PhC  
50 dimensions. Consequently, the mentioned fiber tip E-field sensor [9] utilizes a PhC with an area  
51 of  $19\ \mu\text{m} \times 19\ \mu\text{m}$  which requires, due to its compact size, more complex and expensive systems  
52 such as the FIB system for the precise welding between the PCS and the fiber facet. This  
53 process, due to its serial nature, generally lasts several hours for each sensor head, leading to a  
54 low throughput-integration process remaining still far from massive production procedures. As  
55 an alternative, transfer-printing techniques based on epoxy solutions have been employed for  
56 reducing assembly costs for the welding [13]. However, some constraints remain concerning  
57 the parallelism between the fiber and the photonic structure. A promising technique using a  
58 camera as feedback loop for the optimization of the angle between both surfaces has been  
59 recently reported [14]. Nonetheless, this technique does not take any optical response  
60 information as feedback and, as a consequence, the device can only be tested at the end of the  
61 assembly. In this paper, we demonstrate a cost-efficient assembly method which simplifies the  
62 previously mentioned techniques. Thus, a system based on a stretchable sample holder ensures  
63 the parallelism between both surfaces when applying a vertical force, removing the need of  
64 complex rotational systems. In addition, an active positioning system is also implemented  
65 which allows, by observing the spectrum of the GRs during the assembly, the control of  
66 different parameters such as the beam centering. This novel assembly method can be easily  
67 used for the batch production of different GR-based fiber tip sensors, for a large range of  
68 applications, beyond E-field sensors, like the measurement of physical parameters such as  
69 pressure, refractive index, temperature, etc.

## 70 2. The stretchable membrane concept

71 The fiber-tip E-field sensor developed in this work is sketched in Fig. 1(a,b). An optical fiber  
72 is pigtailed by a ferrule to obtain a surface wide enough to act as a support for 3mm-diameter  
73 photonic chip which represents the sensing part of the device. The chip contains a PhC  
74 structure which is designed on a 700nm-thick membrane of X-cut LN centered to the optical  
75 fiber core.

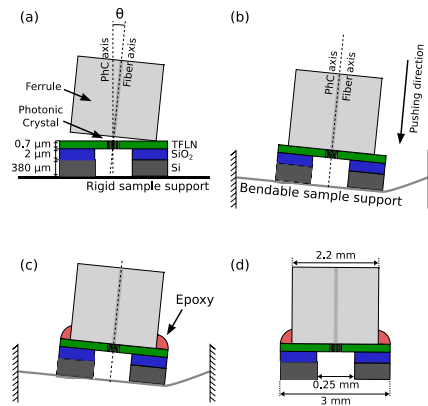


76  
77 Fig. 1. Rendered images of (a) the fully assembled fibered sensor head and (b) sectional cut.

78  
79  
80 The membrane is based on a rigid  $380\ \mu\text{m}$ -thick substrate where only a  $250\ \mu\text{m}$ -diameter circle  
81 located at its center defines the shape of the free-standing membrane. This configuration acts  
82 as a protection crown for the PhC whose sensor performances depend on. The PhC geometry  
83 study and optimization are beyond the scope of this work but explained in detail in [15]. This  
84 thin free-standing PCS can be fabricated from thin film Lithium Niobate-on-insulator (LNOI)  
85 wafers, which are nowadays commercially available thanks to ion slicing techniques [16]. The

86 fabrication of the PCS on the LN suspended membrane has been assessed in a recent work [11],  
 87 demonstrating that highly precise and reproducible spectral performances are achievable.  
 88 However, the existence of an angle  $\theta$  between the fiber facet and the PhC plane may produce a  
 89 serious deterioration of the optical performances, strongly increasing the losses, as  
 90 demonstrated in [9]. Moreover, the macroscopic size of the ferrule makes even more critical  
 91 the control of the parallelism since bigger airgaps, between the PCS and the fiber, are  
 92 introduced when  $\theta > 0$ , as depicted in Fig. 2(a).

93 This non-parallelism angle can be minimized to negligible levels by replacing the rigid sample  
 94 holder by a stretchable membrane. Thus, by applying a force, the membrane is bent, reducing  
 95 the angle error as shown in Fig. 2(b). The applied force may break the 700nm-thick free-  
 96 standing membrane. However, the rigid Si-substrate gives enough robustness to the chip to  
 97 keep its integrity during the process. Once a strong force is applied to the stretchable membrane,  
 98 both surfaces are forced to  $\theta \approx 0$ . Then, the epoxy glue can be applied to weld the ferrule and  
 99 the TFLN which contains the PhC. The applied force on the ferrule towards the chip produces  
 100 the necessary friction to avoid any displacement of the different elements that can be induced  
 101 during the epoxy deposition. The epoxy is hardened during 3 hours at room temperature and  
 102 then the ferrule is pulled back, leading to the assembly sketched in Fig. 2(d). The TFLN is  
 103 directly in contact with the ferrule surface in order to minimize the air- gap between both  
 104 surfaces that might lead to Fabry-Perot effect that can deteriorate the spectral performances. As  
 105 sketched in Fig. 2(d), the chip diameter is slightly higher than the ferrule diameter in order to  
 106 avoid the epoxy to be deposited on the PhC side. However, issues such as aligning both fiber  
 107 and PhC centers or ensuring the optimal optical response are not considered in the standalone  
 108 approach. For this reason, we assist the stretchable membrane approach by an active positioning  
 109 system as further described.



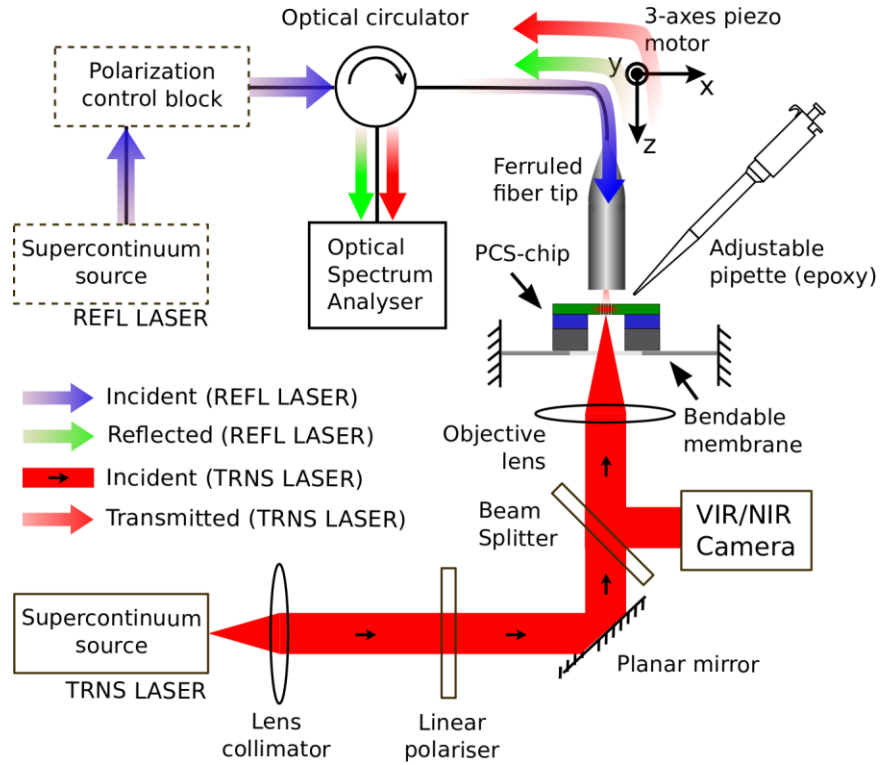
110  
 111 Fig. 2. Assembly of the ferrule-PCS: (a) Fiber approach to the PCS. (b) Pushing the fiber towards the stretchable  
 112 membrane. (c) Epoxy deposition. (d) Fiber-tips sensor release.

113  
 114

### 115 3. The Assembling system

116 The experimental system which fulfills the mentioned requirements is represented in Fig. 3.  
 117 The system employs a supercontinuum source which allows a broadband characterization of  
 118 the transmission spectrum during the assembly. In order to measure the transmitted spectrum,  
 119 the light is firstly polarized in free space as the photonic crystal structure presents different  
 120 Fano resonances as function of the incident polarization. Then, the laser is introduced into an  
 121 inverted microscope which, using an 20x optical objective, injects the light precisely into the  
 122 PhC structure. The alignment process requires the observation of the laser beam with respect

123 to the PhC surface. For this reason, a VIR/NIS is introduced next to the objective to perform a  
 124 first coarse alignment. The sample is, as previously described, located on a stretchable  
 125 membrane.

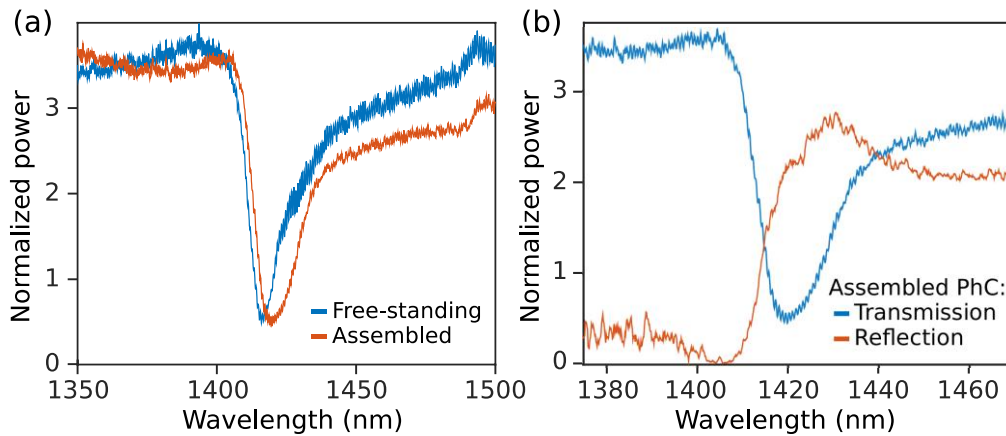


126  
 127

Fig. 3. Experimental system including the dynamic positioning features.

128

129 To avoid any perturbation to the microscope observation, the membrane presents a hole at its  
 130 center. At this point, the ferrule can be approached by employing a 3-axes nano-precision  
 131 piezoelectric to an approximate distance of  $1\mu\text{m}$  away from the PhC. The ferrule can be firstly  
 132 aligned by introducing any VIS/NIR light into the fiber, which transmits through the PhC and  
 133 is captured by the camera. Once aligned, the supercontinuum light is collected by the ferrule  
 134 and monitored by an Optical Spectrum Analyzer (OSA). The transmission at this distance is  
 135 shown by the blue curve in Fig. 4(a), which is normalized with respect to the transmitted  
 136 intensity of the cross polarization



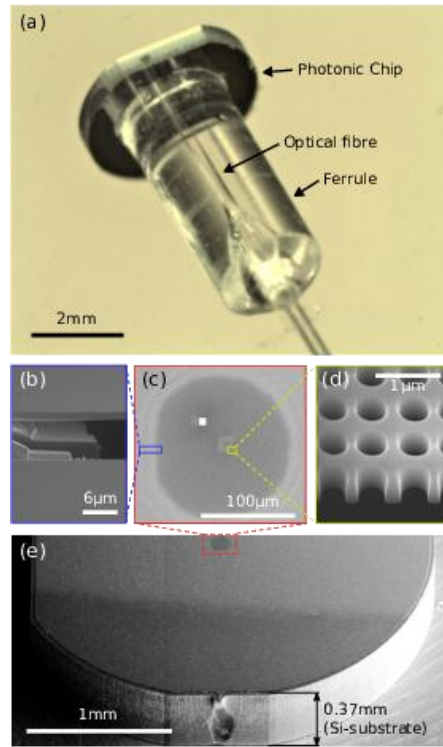
137

138  
139

Fig. 4. (a) Normalized transmission spectra for the free-standing PhC (blue) and for the assembled PCS (red). (b) Comparison of both transmission (blue) and reflection (red) of the assembled PCS.

140  
141  
142  
143  
144  
145  
146  
147  
148  
149  
150  
151  
152  
153  
154  
155  
156  
157  
158  
159

A sub-micrometer precision alignment is performed by optimizing both transmitted power and spectral extinction ratio as a function only of the x and y axes, reducing the degrees of freedom to the translation of x and y. Then, the epoxy is applied with an adjustable pipette. After 1 hour of curing, both transmission and reflection spectra are characterized. The resulting transmission spectrum from the assembly, represented by the red curve in Fig. 4(a), reveals that the spectral slope remains above 0.6dB/nm and does not degrade after the assembly. However, a shift of the resonant wavelength of approximately 5 nm is observed between both curves which can be attributed to the correction of the angle  $\theta$  that changes the light incidence angle into the PhC [17]. In order to obtain the reflection measurement, the supercontinuum source is plugged into the optical circulator. Since the PhC structure is polarization sensitive, the light injected into the ferrule has to be linearly polarized and aligned with the PhC structure. For this reason, a fiber stretching system is installed prior to the optical circulator. A feedback loop for the polarization control can be set by inserting a linear polarization analyzer to the camera, whose axis is aligned with respect to the PhC. The reflection spectrum is represented by the red curve in Fig. 4(b). It can be observed that both curves are complementary. The assembled sensor head is depicted in Fig. 5(a). The photonic chip, the fiber ferrule and the fiber can be clearly distinguished. A SEM picture of the photonic chip, before its assembly onto the ferrule, is shown in Fig. 5(e). As previously described, most of the chip contains the 370 $\mu$ m-thick Si substrate insuring robustness to the sensor head. Despite its circular shape, the chip presents a cut on a side in order to mark the orientation of the PhC and the LN crystallographic axes.



160

161 Fig. 5. (a) Assembled sensor. (b) SEM picture of the sectional cut in a border of the membrane. (c) Top view of the  
 162 free-standing membrane with the PhC structure in the middle. (d) PhC structure on the free-standing membrane. (e)  
 163 SEM image of the chip with the membrane and the PhC

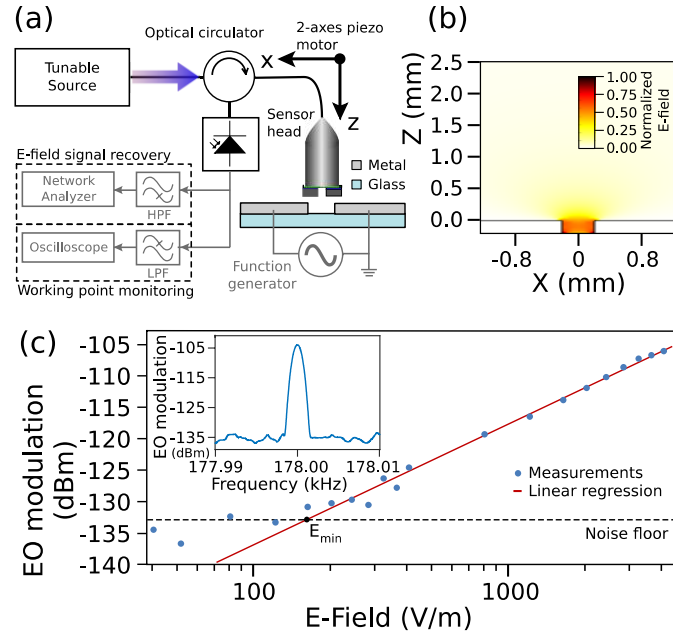
164

165 At the center, the 250  $\mu\text{m}$ -diameter free-standing membrane can be identified due to its darker  
 166 tone, produced by the lower reflection of the SEM electrons induced by its ultra-thin thickness  
 167 in comparison with the non-released areas. The membrane is zoomed in Fig. 5(c), where the  
 168 PhC, whose sectional cut view shown in Fig. 5(d), is located at the center of the membrane. We  
 169 notice that a bright square can be distinguished which corresponds to a drift-correction mark  
 170 employed during the FIB fabrication.  
 171

#### 172 4. Electric field sensing measurements

173 Due to the electro-optic (EO) properties of LN, any applied electric field along the  
 174 crystallographic Z-axis will induce a linear modification of the refractive index with respect to  
 175 the electric field. This modification leads to a shift on the resonant frequency of the Fano  
 176 resonance. Thus, a continuous wave (CW) source has to be employed, whose wavelength is  
 177 located within the Fano spectral slope. The CW light is then modulated by the applied E-field,  
 178 whose modulation strength can be optimized by setting the wavelength at the maximum spectral  
 179 slope of the Fano resonance [11]. In order to evaluate the EO performances, the fibered sensor  
 180 is approached to a PCB (printed circuit board) containing two coplanar transmission lines where  
 181 a time-variant periodic signal is applied as depicted in Fig. 6(a). The tunable source wavelength  
 182 is set at 1576 nm ensuring the optimal EO modulation strength and linearity with respect to the  
 183 applied E-field [9]. Thus, a time-variant periodic signal is applied on the coplanar lines. During  
 184 this experiment, the distance between the PC and the electrodes is approximately 600  $\mu\text{m}$ . In  
 185 addition, the position along the x axis is set to be a at the center of the gap region between the

186 two electrodes since, as determined through electrostatic Finite Element Method (FEM)  
 187 simulations, the maximum E-field value is located in this region as depicted in Fig. 6(b). The  
 188 frequency of the introduced signal is set to 178 kHz due to an optimization of crosstalk  
 189 interferences with other equipment surrounding the system. Moreover, the parameters of the  
 190 network analyzer have been optimized to minimize the noise floor which has been estimated to  
 191 be about  $-132.5$  dBm.



192  
 193  
 194  
 195  
 196

Figure 6: (a) System for the measurement of the E-field sensing. (b) Spatial distribution of the E-field induced by electrodes. (c) Experimental variations of the EO modulation with respect to the applied E-field.

197 The reflected light, which carries the modulation signal, is introduced in a fast photodiode. The  
 198 detected signal is split into high and DC signals by using its respective electrical filters. The  
 199 DC signal carries the information about the reflected power in the absence of any E-field. This  
 200 is useful since other factors such as a temperature change or deformation of the structure may  
 201 add a shift of the resonant wavelength, displacing the optimal operation point. Thus, by  
 202 monitoring the DC signal, the stability of the Fano sensor during its operation can be observed.  
 203 On the other side, the RF signal carries the E-field information. Thus, the signal is collected in  
 204 a network analyzer whose measured bandwidth is represented in the inset of Fig. 6(c). The  
 205 measured modulation frequency ( $f = 178$  kHz) corresponds to the one at the frequency  
 206 synthesizer and it exhibits an EO modulation strength greater than 30 dB with respect to the  
 207 noise floor. The linearity of the device as a function of the E-field strength is assessed by  
 208 varying its value from 40 V/m to 4 kV/m. The resulting EO modulation strength is represented  
 209 in Fig. 6(b) in function of the applied E-field. A linear regression fit demonstrates a notable  
 210 linearity, revealing the absence of non-linear effects strong enough to distort the linear  
 211 operation regime. However, under 250 V/m the regression error increases probably due to the  
 212 noise. Furthermore, the intersection between the noise floor and the linear regression  
 213 determines the minimum detectable field whose value is estimated to approximately 170 V/m.

214 **5. Conclusions**



215 We have developed a procedure for the assembly of PhC structures on the tip of an optical fiber.  
216 The technique presents several advantages with respect to other solutions. First, the stretching  
217 membrane simplifies the alignment procedure and decreases the alignment degrees of freedom  
218 from 5 to 2. Secondly, the technique includes dynamic positioning features allowing the  
219 observation of the photonic crystal optical response during the x-y axis optimization, allowing  
220 a sub-micrometer precision. This solution can be also employed for the integration of LN thin  
221 layers, also as well as for a broad range of materials that provide different functions such as  
222 temperature, pressure and media sensing. Furthermore, the technique requires relatively  
223 inexpensive equipment and may provide high throughput being compatible for batch process  
224 for the future industrialization of lab-on-fiber sensors.

225 **Funding.** This work was supported by the French Directorate General of Armaments (DGA) (ANR-13- ASTR-0019).

226 **Acknowledgments.** We appreciate the technical support of R.-M. Sauvage, L. Robert, S. Bargiel, T. Faure, E. Atie  
227 and R. Zegari. We are also grateful to the RENATECH network and the computing center of Franche-Comte.

228 **Disclosures.** The author declares no conflicts of interest.

## 229 References

- 230 1. J.I. Peterson, G.G. Vurek, "Fiber-optic sensors for biomedical applications," *Science* **224**, 123-127 (1984).
- 231 2. Z. Zhou, J. He, M. Huang, J. He, G. Chen, "Casing pipe damage detection with optical fiber sensors," *Adv.*  
232 *Civil Eng.* **9**, 638967 (2010).
- 233 3. H. Chikh-Bled, K. Chah, A. Gonzalez-Vila, B. Lasri, C. Caucheteur, "Behavior of femtosecond laser-induced  
234 eccentric fiber Bragg gratings at very high temperatures," *Opt. Lett.* **41**, 4048-4051 (2016).
- 235 4. D. Tailaert, "Grating couplers as interface optical fibres and nanophotonic waveguides," PhD dissertation,  
236 Guent University (2004).
- 237 5. C. Nottbohm, A. Turchanin, A. Beyer, R. Stosch, A. Golzhauser, "Mechanically stacked 1-nm-thick carbon  
238 nanosheets, ultrathin layered materials with tunable optical, chemical and electrical properties," *Small* **7**, 874-  
239 883 (2011).
- 240 6. M. Passoni, D. Gerace, L. Carroll, L.C. Andreani, "Grating couplers in silicon-on-insulator: the role of photonic  
241 guided resonances on lineshape and bandwidth," *Appl. Phys. Lett.* **110**, 041107 (2017).
- 242 7. M.L. Brongersma, Y. Lui, S. Fan, "Light management for photovoltaics using high-index nanostructures,"  
243 *Nature Mater.* **13**, 451-460 (2014).
- 244 8. J.R. Piper, S. Fan, "Total absorption in a graphene monolayer in the optical regime by critical coupling with a  
245 photonic crystal guided resonance," *ACS. Photonics* **1**, 347-353 (2014).
- 246 9. V. Calero, M.-A. Suarez, R. Salut, A. Caspar, F. Behague, L. Galtier, G. Gaborit, L. Gilette, F. Baida, N.  
247 Courjal, M.-P. Bernal, "An ultra-side bandwidth, ultra-high spatial resolution and ultra compact electric field  
248 sensor based on Lab-on-fiber technology," *Sci Rep.* **9**, 8058 (2019).
- 249 10. G. Gaborit, J. Dahdah, F. Lecoche, P. Jarrige, Y. Gaeremynck, E. Duraz, L. Duvillaret, "A nonperturbative  
250 electrooptic sensor for in situ electric discharge characterization," *IEEE Trans. Plasma Sci.* **41**, 2851- 2857  
251 (2013).
- 252 11. V. Calero, M.-A. Suarez, R. Salut, B. Robert, A. Caspar, F. Baida, N. Courjal, M.-P. Bernal, "Towards highly  
253 reliable, precise and reproducible fabrication of photonic crystal slabs on lithium niobate," *J. Lightwave*  
254 *Technol.* **37**, 698 (2019).
- 255 12. R. Geiss, J. Brandt, H. Hartung, A. Thinnermann, T. Pertsch, E.B. Kley, F. Schrepel, "Photonic  
256 microstructures in lithium niobate by potassium hydroxide assisted ion beam-enhanced etching," *J. Vac. Sci.*  
257 *Technol. B* **33**, 010601 (2015).
- 258 13. S. Kim, J. Wu, A. Carlson, S. Jin, A. Kovalsky, P. Glass, Z. Liu, N. Ahmed, S. Elgan, W. Chen, P. Ferreira, M.  
259 Sitti, Y. Huang, J. Rogers, "Microstructures elastomeric surfaces with reversible adhesion and examples of  
260 their use in deterministic assembly by transfer printing," *Proc. Natl. Acad. Sci. U.S.A.* **107**, 17095 (2010).
- 261 14. A.Kudryavtsev, G.J.Laurent, C.Clevy, B.Tamadarte, P.Lutz, "Analysis of cad model-based visual tracking for  
262 microassembly using a new block set for matlab/simulink," *International Journal of Optomechatronics* **9**, 295-  
263 309 (2015).
- 264 15. W. Qiu, A. Ndao, H. Lu, M.-P. Bernal, F. Baida, "Guided resonances on lithium niobate for extremely small  
265 electric field detection investigated by accurate sensitivity analysis," *Opt. Express* **24**, 20196-20209 (2016).
- 266 16. <http://www.nanoln.com>
- 267 17. L. Chen, Z. Qiang, H. Yang, H. Pang, Z. Ma, W. Zhou, "Polarization and angular dependent transmission on  
268 transferred nanomembrane Fano filters," *Opt. Express* **17**, 8396-8406 (2009).
- 269