

Correlation between physical properties and electrochromic performances of DC magnetron sputtered  $\alpha$ -WO<sub>x</sub> thin films

K. B. Joël-Igor N'Djoré<sup>a, b</sup>, M. Grafouté<sup>a</sup>, A. Bulou<sup>c</sup> C. Rousselot<sup>\*b</sup>

<sup>a</sup> Université Félix Houphouët Boigny, Laboratoire de Technologie, UFR SSMT 22, BP 258

Abidjan, Côte d'Ivoire

<sup>b</sup> Université de Franche Comté, Institut FEMTO-ST (CNRS/UFC/ENSMM/UTBM),

Département MN2S, BP 71427, F-25211 Montbéliard cedex, France

<sup>c</sup> Le Mans Université, Institut des Molécules et Matériaux du Mans (CNRS UMR6283), F-

72085 Le Mans cedex 09, France

\*Corresponding author: C. Rousselot ([Christophe.rousselot@univ-fcomte.fr](mailto:Christophe.rousselot@univ-fcomte.fr))

4, place Tharradin 25211 Montbéliard Cedex, France, (+33) 381 994 714)

## Abstract

Tungsten oxide thin films were grown on different substrates by the conventional process of direct current magnetron sputtering. A series of amorphous  $\text{WO}_x$  (a- $\text{WO}_x$ ) films, with variable stoichiometry and porosity, were achieved with different total pressure ( $P_t$ ) induced by oxygen flow rate ( $Q(\text{O}_2)$ ) variation. The objective of this study is to determine the couple ( $Q(\text{O}_2)$ ;  $P_t$ ) necessary to achieve optimal electrochromic performances of  $\text{WO}_x$  and also study the correlation between physical properties and electrochromic performances. The  $\text{WO}_x$  films optical transmittance evolved from opaque to transparent with the increase of  $Q(\text{O}_2)$ . A degradation of films refractive index and the increase of their porosity, both mainly due to the increase of  $P_t$ , were observed. The Raman spectroscopy indicated the presence of different W valence states in the as deposited a- $\text{WO}_x$  films. They all contain  $\text{W}^{4+}$ , but  $\text{W}^{6+}$  were found only into transparent or semi-transparent films. A weak peak of  $\text{W}^{5+}$  was observed into the semi-transparent film. Electrical properties studied by Metal-Oxide-Semiconductor structures showed a decrease of the oxide capacitance when  $Q(\text{O}_2)$  and  $P_t$  increase. The  $\text{WO}_x$  transparent and optically reversible films, deposited with lower values of ( $Q(\text{O}_2)$ ;  $P_t$ ), i.e.,  $Q(\text{O}_2) = 4.5$  sccm and 6 sccm, and  $P_t = 1.13$  Pa and 1.37 Pa, respectively, exhibited the highest values of  $\text{WO}_x/\text{Si}$  interface traps average density and positive ions density ( $N_{ss}$ ). The positive  $N_{ss}$  reflects the oxygen deficiency amount in the films. These films were resistive to the discoloration process but showed the highest coloration efficiency  $CE_{632\text{nm}} = 44.4$   $\text{cm}^2/\text{C}$  and 42.2, respectively. Their coloration times were  $t_c = 19$  s and 22 s, and their discoloration times were  $t_d = 180$  s and 75 s, respectively. They exhibited the highest optical modulation close to  $\Delta T_{632\text{nm}} = 76$  %.

## Key words

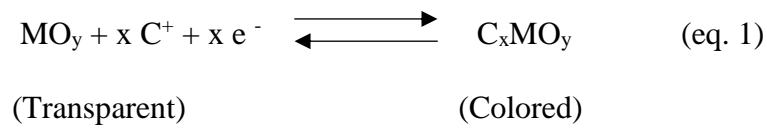
- $\text{WO}_x$  thin films
- Direct Current magnetron sputtering
- Optical properties
- Electrical properties
- Metal-Oxide-Semiconductor structures
- Raman spectroscopy
- Electrochromic performances

## Highlights

- $\text{WO}_x$  thin films ( $0 < x \leq 3$ ) deposited by DC reactive magnetron sputtering
- Variation of the porosity and the stoichiometry of  $\text{WO}_x$  films by total pressure and  $Q(\text{O}_2)$
- Optical, electrochromic and electrical properties studied and correlated with  $Q(\text{O}_2)$
- Electrochromic properties are correlated to electrical parameters such as  $D_{it}$  and  $N_{ss}$
- Raman study reveals different valence states,  $\text{W}^{4+}$ ,  $\text{W}^{5+}$  and  $\text{W}^{6+}$  depending of  $Q(\text{O}_2)$

## 1 Introduction

Electrochromic materials are able to change their optical properties under the action of an electrical voltage. This change of optical properties when a voltage is applied across electrochromic materials is called electrochromism [1]. Electrochromism for cathodic coloring materials, often transition metal oxides, is based on a double injection/extraction of electrons and cations in/from their structure according to the following redox equation [2]:



where M is a transition metal (W, Ta, Ti, etc.), C is a cation ( $\text{Na}^+$ ,  $\text{Li}^+$ ,  $\text{H}^+$ , etc.) and x is the fractional number of inserted ions and electrons inside the  $\text{MO}_y$  film.

Electrochromism is known to be existed in several types of transition metals oxide [2]. Among them, tungsten trioxide ( $\text{WO}_3$ ) is by far the most extensively studied inorganic electrochromic material because of its cubic “perovskite-like” structure [2,3]. The free space inside the  $\text{WO}_3$  cubic structure is considerable and provide a large number of interstitial sites where cations can be inserted. The coloration process of  $\text{WO}_3$  is described by different models which are sometimes contradictory [4]. One of the most accepted models to explain  $\text{WO}_3$  coloration is the intervalence charge transfer transition between  $\text{W}^{5+}$  and  $\text{W}^{6+}$  sites after the insertion of electrons and cations [5,6].  $\text{WO}_3$  coloration is also explained by small polaron absorption [7]. As  $\text{WO}_3$  optical transmittance is reversible, the disinsertion of electrons and cations allows this material to retrieve its initial transparent state. Due to its optical change, the  $\text{WO}_3$  thin film can be integrated in devices that modulate their transmittance, reflectance or absorptance [8,9]. Tungsten oxide thin films are generally used as an active layer in several

optoelectronics devices such as architectural glazing, display screen, sun roof, solar glass and smart windows [1,8,10]. A large variety of deposition techniques like thermal evaporation [11], spray pyrolysis [12], chemical vapor deposition [13] and magnetron sputtering [14] have been utilized to obtain  $\text{WO}_3$  thin films.

According to literature [15,16], amorphous  $\text{WO}_3$  thin films with relatively porous microstructure and larger thickness exhibit optimal electrochromic (EC) performances, good optical reversibility and smallest coloration/bleaching time. Porous films can be obtained with high total pressure during deposition and amorphous films are generally deposited at room temperature [15]. For Yu et al. [17] and Lee et al. [18], amorphous and sub-stoichiometric  $\text{WO}_{3-y}$  films exhibit better optical modulation and coloration efficiency. So, the attractive EC performances are due to the synergistic effects of the porosity and oxygen vacancy (or deficiency) in the  $\text{WO}_x$  [19]. Therefore, the choice of the deposition technique, deposition conditions and parameters are of a great importance to achieve an electrochromic material with optimal performances.

The oxygen vacancy defects are supposed to increase  $\text{WO}_x$  conductivity and influence their electrochromic performances [20]. But since O-rich  $\text{WO}_x$  films are electrically resistive, it is difficult to measure their conductivity or resistivity with a classic 4-probes device. The oxide films electrical properties in this study will be studied from Metal-Oxide-Semiconductor (MOS) structures. Indeed, the electrical characteristics of the MOS structures can help us to assess defects in oxides films [21]. In  $\text{WO}_x$  films, these defects accessible from a MOS structure are of a great importance in the understanding of their electrochromic behavior. The  $\text{WO}_x$  electrical properties studied by MOS structures is not widespread in the literature and their correlation with electrochromic performances is not well-known.

In this study, we sputter-deposited  $\text{WO}_x$  films by Direct Current (DC) magnetron sputtering with different oxygen flow rates ( $Q(\text{O}_2)$ ) and total pressures ( $P_t$ ) in order to modulate their stoichiometry and porosity, necessary for the optimization of their electrochromic performances. We also examine the influence of  $Q(\text{O}_2)$  and  $P_t$  on optical and electrical properties and the correlation of these properties with the electrochromic performance of  $\text{WO}_x$  films. We will emphasize on the correlation between electrical properties and electrochromic performances of  $\text{WO}_x$ . The ions density ( $N_{ss}$ ) and the interface traps density distribution ( $D_{it}$ ) at the interface  $\text{WO}_x/\text{Si}$ , constitute the main electrical parameters, which will be correlated to the electrochromic performance, such as the coloration efficiency and the kinetics of ions insertion/extraction in/from  $\text{WO}_x$  structure.

## **2 Experimental procedures**

### **2.1 $\text{WO}_x$ films deposition**

Tungsten oxide films were deposited by DC magnetron sputtering of a W target (purity 99.9 % and diameter 2") in a gas mixture of Argon and oxygen. A constant current of 0.1 A was applied to the W target. Reactive sputtering experiments were performed using a vacuum reactor (Alliance Concept AC 450) of 70 L. Before introducing constant Ar and  $\text{O}_2$  gas flow rates in the reactor, a base pressure of  $10^{-4}$  Pa was obtained with a turbomolecular pump backed with a mechanical one. The pumping speed during deposition process was  $10 \text{ L s}^{-1}$ . Amorphous films were deposited on glass and  $\text{SnO}_2$  substrates at room temperature using different total pressures and different constant oxygen flow rates for an Ar pressure maintained at 0.41 Pa. The substrates were cleaned with ethanol and dried at room temperature before each sputtering process. The glass/ $\text{SnO}_2$  substrates were ultrasonically cleaned during 8 minutes to remove eventual soils, then rinsed in distilled water and finally dried at room temperature. For each

deposition, all deposit parameters were recorded using the deposit monitoring software. The WO<sub>x</sub> thin films thickness deposited onto glass substrates were measured using a Dektak 3030 profilometer. In table 1, we have listed the oxygen flow rates, total pressures, oxygen stoichiometries and all the films thickness deposited on glass or glass/SnO<sub>2</sub> substrates.

## 2.2 Characterizations

The content in O<sub>2</sub> and W atom in deposited films was obtained by Energy Dispersive X-Ray Spectroscopy (EDS). For this purpose, we used an OXFORD detector in the FEG JEOL 7610 F microscope. To achieve the quantitative analysis of O<sub>2</sub> and W in the films, K<sub>α</sub>-line (0.525 eV) and M<sub>α</sub>-line (1.775 eV) signals were used respectively. The acceleration voltage and the current for the EDS measurements were 5 kV and 100 nA. The atomic ratios of oxygen and tungsten present in the deposited films, are shown in Table 1.

The SEM (Scanning Electron Microscope) measurements were carried out in order to observe the morphology of the deposited films. Before the SEM analyses, the samples were covered with a thin metallic layer of chromium (20 nm of Cr) in order to ensure the electrical conductivity of the samples and to avoid any drift due to charging effects during the FIB cutting process (Focused Ion Beam). This Cr deposit was carried out by magnetron sputtering in a Leica ACE 600 metallizer. The samples were then installed in a dual-beam SEM/FIB station, Helios Nano lab 600i from FEI. The FIB column generates a focused Ga<sup>+</sup> ion beam in order to obtain a cross-section of the sample.

Structural X-Ray Diffraction (XRD) analyzes reveal no peak characteristic of a crystallized WO<sub>3</sub> structure. It can therefore be confirmed that all the a-WO<sub>x</sub> films deposited at room temperature are amorphous (see the insert of Figure 1.b).



Raman spectroscopy study was carried out on a-WO<sub>x</sub> thin films of about 500 nm of thickness deposited onto glass substrates in order to determine the chemical bonding configuration and W valence states. The Raman scattering spectra were collected with a T64000 multichannel spectrometer (Horiba-Jobin-Yvon) using a Si-based CCD detector cooled to – 133°C (liquid nitrogen-cooled). The experiments were performed under optical microscope, (Olympus BX41), in the backscattering geometry, with a X50 LF objective (numerical aperture 0.5) and using as excitation the 514.5 nm wavelength radiation (Ar/Kr Coherent Spectrum laser) with power limited to 0.6 mW on the sample. The spectra were collected in a broad wavenumber range, from 120 cm<sup>-1</sup> to 3500 cm<sup>-1</sup> (single monochromator and 600 tr/mm grating/1.8 cm<sup>-1</sup> instrument resolution), and close to the low wavenumbers, from 20 cm<sup>-1</sup> to 1500 cm<sup>-1</sup> (triple monochromator and 1800 tr/mm grating/ 0.7 cm<sup>-1</sup> instrument resolution). They were calibrated with the 520.2 cm<sup>-1</sup> line of a silicon wafer and corrected from glass substrate signal. The Raman spectra wavenumbers, intensities, and widths were then fitted by Labspec5 software using a combined Lorentzian-Gaussian band shape.

Films deposited on glass substrates were characterized by optical transmission spectroscopy with normal incidence, measured with a Perkin Elmer lambda 950 spectrophotometer, in the wavelength range from 200 nm to 1,100 nm, in order to draw their transmission spectrum. For each spectrum, the refractive index of each WO<sub>x</sub> films is determined from the interference fringes, using the Swanepoel method [22,23].

For electrical study, films with only about 200 nm were deposited onto n-type Si wafers (100) with a resistivity of 1-20 Ω.cm and a carrier concentration of 10<sup>14</sup>-10<sup>15</sup> cm<sup>-3</sup>. The electrical measurements were performed with a mercury probe capacitance meter. The electrical procedure measurement with this capacitance meter, is widely detailed in our previous paper [24]. Electrical characteristics studied in this work were the oxide capacitance (C<sub>ox</sub>), the

hysteresis of the capacitance-voltage curve: C-V ( $\Delta V$ ), the interface traps density distribution ( $D_{it}$ ) [25,26] at  $WO_x/Si$  interface and ions density ( $N_{ss}$ ). The  $D_{it}$  and  $N_{ss}$  are calculated from the indications given by N'Djoré et al [24].  $C_{ox}$  is measured in accumulation regime. **Hysteresis is determined by applying a back-and-forth gate voltage across the MOS structures. The voltage evolves from 0 V toward negative voltages, then reverse to positive voltages and end at 0 V.** Among the methods that exist to quantify defects in the volume of oxide films, we used the hysteresis loop of C-V curves by applying a back-and-forth gate voltage ( $V_g$ ) with a frequency of 1 MHz across MOS structures. The mid-height width of the hysteresis loop of C-V curve, appointed  $\Delta V$ , is representative of defects in the volume of  $WO_x$  films [21].

Thin films of  $WO_x$  with about 500 nm [27] of thickness were deposited onto glass/ $SnO_2$  substrates for electrochromic study. The  $SnO_2$  has a square resistance of  $26 \Omega/\square$ . After deposition of  $WO_x$  films, electrochemical analyses were performed with a three electrodes electrochemical cell. The electrochemical cell is composed of a working electrode (glass/ $SnO_2/WO_x$ ) with a contact surface of  $2 \text{ cm} \times 1.5 \text{ cm} = 3 \text{ cm}^2$ , a saturated reference electrode  $Hg/Hg_2Cl_2/KCl$ , noted SCE (Saturated Calomel Electrode) and a counter electrode which consist of a platinum grid. The three electrodes were immersed in a 1.0 M solution of lithium perchlorate ( $LiClO_4$ ), dissolved in propylene carbonate. Electrodes were connected to a 33V/2A potentiostat fitted with an intelligent programmable interface (IMT 101) and coupled to the VoltaMaster 2 software for data recording and processing during the cyclic voltammetry and chronoamperometry measurements. Voltammetry measurements were conducted between - 1 V and + 1 V vs SCE with the scan rate of 20 mV/s in order to observe the phenomena of intercalation and desintercalation of carriers ( $Li^+$ ,  $e^-$ ) through the current density-voltage (J-V) curves. Chronoamperometry measurements were done at a fix potential of - 1 V for films coloration (reduction reaction). Then, after coloration, a fix voltage of + 1 V was applied to

achieve films discoloration (oxidation reaction). Chronoamperometry measurements permitted us to evaluate films response time, optical modulation and coloration efficiency.

### 3 Results and discussions

In table 1, we have listed the oxygen flow rates, total pressures, oxygen stoichiometries and all the films thickness deposited on glass or glass/SnO<sub>2</sub> substrates. The same films (table 1), but with thickness of about 200 nm, were deposited onto silicon and in the same conditions for electrical measurements.

#### 3.1 Optical properties and porosity

The Figure 1.a shows the WO<sub>x</sub> optical transmission versus the wavelength for various oxygen flow rates between 1.2 and 14 sccm and (i.e., total pressure 0.45 and 2.60 Pa). The WO<sub>x</sub> thin films average optical transmittance calculated in the visible region between 400 nm and 800 nm versus oxygen flow rate is shown also in Figure 1.b.

For  $Q(O_2) = 1.2$  sccm (WO<sub>1.92</sub>), the WO<sub>x</sub> sample obtained is opaque. Its average optical transmittance is close to zero in the visible region. When  $Q(O_2) \geq 1.65$  sccm, films evolve from semi-transparent (bluish) to transparent with appearance of interference fringes in the visible region. The bluish film deposited with  $Q(O_2) = 1.65$  sccm (WO<sub>2.56</sub>), presents an average optical transmittance of 32% in the visible range. The blue aspect of the film is due to an under-stoichiometric, with considerable oxygen vacancies associated to W<sup>4+</sup> and W<sup>5+</sup> [28]. This under-stoichiometric permits the absorption of the photon which induce an electronic transfer from between the O valence band to the W conduction band, at the origin of the bluing. In addition, the bluing can also be attributed to an electronic inter-valence transition between W<sup>6+</sup> and W<sup>5+</sup>. Finally, all the WO<sub>x</sub> films deposited with  $Q(O_2) \geq 1.95$  sccm have an average optical transmittance of about 80 % and are transparent with interference fringes. Indeed, the films

prepared with  $Q(O_2) \geq 3$  sccm become increasingly oxygen-rich to reach the stoichiometry  $WO_3$  at  $Q(O_2) = 14$  sccm (Table 1) with a reinforcement of covalent bonds W-O. This reinforcement of W-O bonds may disadvantage the inter-atomic electronic transfer and also the electronic inter-valence transition between  $W^{6+}$ ,  $W^{5+}$  and  $W^{4+}$ , reason why these deposited films are transparent.

According to Figures (1.a and 1.b) and table 1, we can assume that the transparency of films deposited between  $3 \text{ sccm} \leq Q(O_2) \leq 14 \text{ sccm}$  is probably due to their enrichment in oxygen and the predominance of  $W^{6+}$  as observed by Li et al.[24,29]. As a consequence, these transparent films could be potential candidates for electrochromic studies.

The refractive index ( $n$ ) of  $WO_x$  thin films was determined from the refractive index of the glass substrate and the successive optical transmittance maxima and minima of optical transmittance interference fringes (Figure 1.a) by using Swanepoel's method [22,23]. Changes in oxygen flow rate and total pressure allow the formation of films with different porosities and result in different refractive indices. We choose to represent in Figure 2, the effect of total pressure on the evolution of the refractive index, only for transparent  $WO_x$  thin films. The refractive index is represented versus the wavelength with  $425 \text{ nm} \leq \lambda \leq 800 \text{ nm}$ .

Figure 2 shows a progressive decay of  $WO_x$  thin film refractive index when the total pressure in the deposition reactor increases. This result is in accordance with Sun et al.[15] study. In addition, we observe in this Figure 2 a systematic decrease of the refractive index of each  $WO_x$  thin film when the wavelength evolves from 425 nm to 800 nm. This refractive index dispersion with respect to the wavelength in visible region, is well known for transparent films and is characterized by the following Cauchy equation [30] :

$$n(\lambda) = A + B/\lambda^2 + C/\lambda^4 + \dots \quad (\text{eq.4})$$

Where  $n(\lambda)$  is the refractive index for a given wavelength,  $\lambda$  is the wavelength, and A, B, C are coefficients which are characteristics of the deposited films.

According to literature, the refractive index of transparent  $\text{WO}_x$  thin films depends on their density (or porosity) [15,31,32]. Since the porosity is influenced by oxygen stoichiometry and the total pressure [15], we have represented the evolution of both average refractive index and average porosity of films, as a function of total pressure and oxygen flow rate in Figure 3. The average refractive index and average porosity were calculated in visible region, precisely for  $425 \text{ nm} \leq \lambda \leq 800 \text{ nm}$ . First of all, we calculated the film relative density  $P(\lambda)$  with the Lorentz-Lorentz equation [15,32] :

$$P(\lambda) = \frac{\rho_f}{\rho_b} = \frac{n_f^2(\lambda)-1}{n_f^2(\lambda)+2} \times \frac{n_b^2+2}{n_b^2-1} \quad (\text{eq.5})$$

where  $\rho_f$  and  $\rho_b$  are respectively the density of  $\text{WO}_x$  thin films and  $\text{WO}_3$  bulk,  $n_f$  and  $n_b$  are the refractive index of  $\text{WO}_x$  thin films and  $\text{WO}_3$  bulk. For  $\text{WO}_3$  bulk, the average value of the refractive index is 2.5 [33]. The films porosity  $\pi(\lambda)$  is then obtained by the following expression [34]:

$$\pi(\lambda) = 1 - P(\lambda) \quad (\text{eq.6})$$

According to Figure 3, when total pressure and oxygen flow rate increase, the evolution of the average refractive index and of the average porosity of  $\text{WO}_x$  are inversed. The dense films show higher refractive index compared to porous films. The average refractive index of the  $\text{WO}_x$  films decreases from 2.13 to 2.03 while average porosity increases from ~ 15 % to ~ 20 % for respectively  $P_t = 0.61 \text{ Pa}$  and  $2.6 \text{ Pa}$  and  $Q(\text{O}_2) = 1.95 \text{ sccm}$  and  $14 \text{ sccm}$ . Consequently, we can say that the decay of the average refractive index is due to the increase

of films average porosity, the latter is induced by the increase of total pressure. Some authors already attributed the decrease of  $\text{WO}_3$  [34],  $\text{Ta}_2\text{O}_5$  [35] and  $\text{ZrO}_2$  [36] films refractive index to the increase of their porosity, thus confirming our results. We can also notice that the average refractive index values of  $\text{WO}_x$  thin films presented in Figure 3 are lower than that of the bulk ( $n_b = 2.5$ ). This confirm that  $\text{WO}_x$  films are not dense compared to the bulk and contain porosity. Let us remember that the increase of the porosity is necessary to promote the electrochromic devices performances [37].

### **3.2 Morphology**

SEM observations (Figure 4 ) were performed on cross sections of the films deposited on silicon substrates with oxygen flow rates of 1.2 sccm ( $\text{WO}_{1.92}$ ), 4.5 sccm ( $\text{WO}_{2.85}$ ) and 14 sccm ( $\text{WO}_3$ ). The SEM observations reveal that the patterns in the  $\text{WO}_x$  films sizes seem to decrease with increasing oxygen flow rate and total pressure. In the literature, some authors [29,38] attribute this evolution of the size to the improvement in the stoichiometry of the film. Indeed, when the oxygen flow increases the film tends towards a more stoichiometric  $\text{WO}_3$  formation. The trend of XRD peak width with changing the oxygen flow rate (and total pressure also) could be due to the progressive evolution of films stoichiometry and the modification of their microstructure (porosity). Indeed, the film deposited with  $Q(\text{O}_2) = 14$  sccm exhibited a broad peak whereas the others films exhibited narrow peaks. The larger XRD peak of  $Q(\text{O}_2) = 14$  sccm, stoichiometry close to pure  $\text{WO}_3$ , may be also due to the change in microstructure such as change in grain size as shown by SEM images and the increase of porosity. The others films are under stoichiometric and less porous.

### **3.3 Raman spectra**

The Raman scattering spectra of four a-WO<sub>x</sub> thin films in the range 20-1500 cm<sup>-1</sup> are shown in Figure 5 for various oxygen flow rates between 1.20 and 14 sccm, together with the spectrum of crystallized (monoclinic) WO<sub>3</sub>. All the spectra of the films exhibit broad bands structures characteristic of the amorphous state, and there is no trace of any line from crystallized WO<sub>3</sub>, in agreement with the X-ray diffraction results. The samples prepared with the highest values of Q(O<sub>2</sub>), from 14 sccm to 1.95 sccm, exhibit very similar profiles characterized by four main bands around 960 cm<sup>-1</sup>, 800 cm<sup>-1</sup>, 280 cm<sup>-1</sup> and 50 cm<sup>-1</sup>. On the other hand, the spectra are significantly different for Q(O<sub>2</sub>) = 1.65 sccm where structures appear in the range 200-480 cm<sup>-1</sup>, and drastically different for Q(O<sub>2</sub>) = 1.20 sccm. No other signal is observed up to 3800 cm<sup>-1</sup> for any of these samples.

To interpret these spectra, we refer to the Raman signals of WO<sub>x</sub> thin films with variable oxygen content reported in literature. Signals in the range 950 cm<sup>-1</sup> - 965 cm<sup>-1</sup> are assigned to the stretching mode of the W<sup>6+</sup>=O bond of terminal (non bridging) oxygen, often observed for amorphous a-WO<sub>3</sub> [39–41]. One can find also an intense broad band in the range 600-900 cm<sup>-1</sup> attributed to O-W<sup>6+</sup>-O stretching vibrations [39–41], as deduced from comparison with crystallized WO<sub>3</sub> (monoclinic) where the most intense Raman lines are found at 808 and 714 cm<sup>-1</sup> [42,43] (Figure 5); a weaker intensity broad band around 280 cm<sup>-1</sup> is attributed to O-W<sup>6+</sup>-O bending vibrations [43–47,47,48]. In addition, Lee et al. [45–47,49] showed that bands centered around 330 and 450 cm<sup>-1</sup> can be attributed to O-W<sup>5+</sup>-O and W<sup>5+</sup>=O vibrations, respectively, and that a band centered near 220 cm<sup>-1</sup> could be associated with W<sup>4+</sup>-O vibrations. Bands in the 1,100-1,500 cm<sup>-1</sup> range and around 3,600 cm<sup>-1</sup> are usually attributed to bending vibration of the W-OH bonds and water [50].

The vibrational characteristics are deduced from fitting of the Raman spectra in the spectral range 120-500 cm<sup>-1</sup> and 500-1,100 cm<sup>-1</sup> (Figure 5). Figure 6 shows the refinements for the samples deposited with Q(O<sub>2</sub>) = 1.20, 1.65, and 1.95 sccm that illustrate the three different kinds of spectra. The very narrow peak at 235 cm<sup>-1</sup> observed in all spectra is a contamination

by plasma line. The vibrational parameters (wavenumbers, spectral widths) of the components are summarized in Table 2 with the attribution proposed as deduced from literature. All the components exhibit gaussian profiles except the most intense band around  $800\text{ cm}^{-1}$  that is asymmetric; this signal appears perfectly fitted by two gaussian lines, as already pointed out [41], centered at  $671\text{ cm}^{-1}$  and  $782\text{ cm}^{-1}$  (Figure 6). For all the samples prepared with  $Q(\text{O}_2) \geq 1.95\text{ sccm}$ , this profile remains strictly unchanged (the two components having the same wavenumbers, the same widths, and the same relative intensities). This band, attributed to O- $\text{W}^{6+}$ -O stretching is used as reference. The weak and very broad band centered at  $278\text{ cm}^{-1}$  also keep the same spectral characteristics and strictly the same integrate intensity ratio with respect to the  $671\text{-}782\text{ cm}^{-1}$  band for any sample with  $Q(\text{O}_2) \geq 1.95\text{ sccm}$  (Figure 7a), which supports its attribution to the O- $\text{W}^{6+}$ -O bending [43–47,47,48]. This witnesses for the fact that the main part of the material (bulk) still is composed of the same  $\text{W}^{(6+)}\text{O}_6$  (amorphous) network (the Raman signal being characteristic of the local arrangement). The highest frequency band, centered at  $961\text{ cm}^{-1}$  (with a very small shoulder around  $1,010\text{-}1,020\text{ cm}^{-1}$ ), is attributed to  $\text{W}^{6+}=\text{O}$  stretching vibrations [39], the weak shoulder already reported in this frequency range, being attributed to the asymmetric stretching [51,52].

When  $Q(\text{O}_2) = 1.65\text{ sccm}$ , the weak broad signal centered at  $278\text{ cm}^{-1}$  is replaced by a more intense and structured band unambiguously composed of mainly four lines with about the same magnitudes (Figure 6). The corresponding wavenumbers agree with published works, and these components can be unambiguously attributed (Table 2) to the  $\text{W}^{+4}\text{-O}$  ( $221\text{ cm}^{-1}$ ), O- $\text{W}^{5+}$ -O ( $320\text{ cm}^{-1}$ ), and  $\text{W}^{5+}=\text{O}$  vibrations; in this latter case however, the band at  $430\text{ cm}^{-1}$  reported in literature appears in the present case as composed of two lines at  $458\text{ cm}^{-1}$  and  $423\text{ cm}^{-1}$ . The sample obtained with  $Q(\text{O}_2) = 1.20\text{ sccm}$  exhibits the same lines plus an intense and narrow peak at  $268\text{ cm}^{-1}$ . In these two samples, in the high frequency range the bands  $671\text{-}782$  and  $961\text{ cm}^{-1}$  are less resolved for  $Q(\text{O}_2) = 1.65\text{ sccm}$ , and they totally merge and strongly decrease for



$Q(O_2) = 1.20$  sccm (Figure 6). The spectra are described by an additional signal between these bands, at  $922\text{ cm}^{-1}$  and  $885\text{ cm}^{-1}$  for  $Q(O_2) = 1.65$  and  $1.20$  sccm, respectively.

As pointed out, all the films deposited with  $Q(O_2) \geq 1.95$  sccm exhibit the same lines with the same profiles. On the other hand, the ratio of intensities of the  $961\text{ cm}^{-1}$  line with respect to the  $671\text{-}782$  band (i.e.  $W^{6+}=O / O\text{-}W^{6+}\text{-}O$ ) continuously decreases (Figure 7a) with decreasing  $Q(O_2)$ , and even linearly decreases as a function of the EDS O/W ratio (Figure 7b). It can be seen also (Figures 7) that the lines at  $320\text{ cm}^{-1}$ ,  $423\text{-}458\text{ cm}^{-1}$  characteristic of  $W^{5+}$ -oxygen vibration, remain proportional to the (remaining) band  $671\text{-}782\text{ cm}^{-1}$  characteristic of  $O\text{-}W^{6+}\text{-}O$  stretchings; this suggests that the  $W^{5+}$  are connected to the  $O\text{-}W^{6+}\text{-}O$  network clusters rather than to the  $WO_2$  network. On the other hand, the line characteristic of  $W^{4+}$  at  $221\text{ cm}^{-1}$ , strongly grows from  $Q(O_2) = 1.65$  sccm to  $1.20$  sccm, in agreement with the fact that the sample become close to the  $WO_2$  stoichiometry. Figures 7 also suggest that for the highest values of  $Q(O_2)$ , this signal is too weak with respect to the  $278\text{ cm}^{-1}$  line to be detected. It can be pointed out that even in this sample where the composition is close to  $WO_2$ , with a majority of  $W^{4+}$  (and still a little  $W^{5+}$  and  $W^{6+}$ ), the whole Raman spectrum is weak (with almost the same magnitude as the contamination by low wavenumber rotational lines of air in the near vicinity of the surface- Figure 5). This is at least partly due to the high reflectivity of the surface (as deduced from the transmittance shown fig. 1), but also to a weak Raman scattering efficiency, presumably due to a somewhat different structural arrangement; this is also supported by the fact that the broad signal around  $50\text{ cm}^{-1}$  always observed for high  $Q(O_2)$  samples (Figure 5), corresponding to narrow lines in crystallized  $WO_3$ , slightly decreases for  $1.65$  sccm, and totally vanishes for  $1.20$  sccm samples. The fairly sharp signal at  $268\text{ cm}^{-1}$  and the line at  $885\text{ cm}^{-1}$  could be characteristic of this arrangement. Incidentally, this spectrum looks like the one of crystallized  $WO_2$  [53,54], as well as the Raman spectra of amorphous  $WO_3$  match the spectrum of crystallized  $WO_3$  (Figure 5), as already pointed out [41,45].

These results confirm the presence of a majority of  $W^{4+}$  ions in the film of composition  $WO_{1.92}$  deposited at  $Q(O_2) = 1.20$  sccm and a very low oxygen deficiency (compared to stoichiometric  $WO_2$ ), reason why it appeared opaque after deposition (Figure 1). Then for  $Q(O_2) = 1.65$  sccm, the film contains  $W^{5+}$  in addition to the  $W^{4+}$ , and  $W^{6+}$  valence states, with an intermediate stoichiometry between  $WO_2$  and  $WO_3$ . The presence of  $W^{5+}$  ions in this film seems to be the cause of its bluish color (Figure 1). There is a majority of  $W^{6+}$  ions, which is responsible for their transparency (Figure 1).

### 3.4 Electrical properties

The electrical parameters presented in this section such as  $C_{ox}$ ,  $N_{ss}$ ,  $D_{it-average}$  and the  $\Delta V$  are determined from capacitance-voltage (C-V) curves obtained during a previous study [24].

#### 3.4.1 $WO_x$ capacitance

The evolution of the  $C_{ox}$  of  $WO_x$  films versus oxygen flow rate is showed in Table 3. According to this Table 3, the  $WO_x$  thin films  $C_{ox}$  is influenced by the oxygen flow rate. The capacitance globally decreases from  $C_{ox} = 8,88 \times 10^{-10}$  F to  $4.01 \times 10^{-10}$  F when  $Q(O_2)$  increases from 1.2 sccm to 14 sccm. The  $C_{ox}$  values of  $WO_x$  films can be compared to the values obtained by Tutov [55]. We can observe that the O-rich films present lower values of  $C_{ox}$  compared to the film synthesized with 1.2 sccm. The decay of the  $C_{ox}$  when  $Q(O_2)$  increases may be mainly due to the resistive nature of O-rich films [56]. We must precise that total pressure and porosity may influence the evolution of the  $C_{ox}$ , but we still do not know how.

#### 3.4.2 Defects and ions density in $WO_x$ films

Defects result from the growth conditions of the oxide. In this work, we will focus on  $\text{WO}_x$  intrinsic defects which are due to the arrangement of oxygen and tungsten atoms in the films. Intrinsic defects can appear as interstitials, substitutions, or as oxygen vacancies associated with contracted, broken or dangling bonds [21,26]. These defects are present in the volume of  $\text{WO}_x$ , close to the interface or at the interface  $\text{WO}_x/\text{Si}$ . A defect can be electrically active and can behave like a trapping site when it captures a charge carrier or like a site of recombination when it can exchange a charge carrier with the oxide valence or conduction bands [57,58].

Various methods exist to quantify defects in oxide films. In our study, we used the hysteresis loop of C-V curves by applying a back-and-forth gate voltage ( $V_g$ ) with a frequency of 1 MHz across MOS structures. The hysteresis is a common phenomenon in the dielectrics (oxides) and is considered to be due to charge trapping [21,59]. The mid-height width  $\Delta V$  of the hysteresis loop is representative of defects in the volume of  $\text{WO}_x$  films which are not in contact with the substrate. More  $\Delta V$  amplitude is important, more defects in the volume of  $\text{WO}_x$  films are important. Also, we determine the  $\text{WO}_x/\text{Si}$  interface traps or states density distribution ( $D_{it}$ ) by Terman method [25,26,60]. The interface trap is a defect that lies precisely at the oxide/semiconductor interface, and can communicate electrically with the underlying semiconductor [21]. For Terman, at an oxidized semiconductor surface, the transition from ordered semiconductor crystal lattice to the oxide layer represents a major discontinuity which will result in dangling or unsaturated bonds [26]. The interface traps are similar to defects found in the volume of the oxide.

As seen, oxygen vacancies in  $\text{WO}_x$  films are considered as defects measurable taking into account in the measurements of  $D_{it}$  and also  $\Delta V$ . When the oxygen flow rate increases, the evolution of these oxygen vacancies can be reflected by the evolution of the  $N_{ss}$  in the  $\text{WO}_x$

films. The  $N_{ss}$ ,  $\Delta V$ , and  $D_{it-average}$  for all  $WO_x$  films are shown in Table 3. As an example of determination of  $\Delta V$ , we show in Figure 9 the hysteresis loop of films deposited with  $Q(O_2) = 8$  sccm.

According to the data of Table 3, a negligible hysteresis loop and the lowest  $D_{it-average}$  are observed for the sample elaborated with 1.2 sccm. The result reflects a good interface  $WO_x/Si$  quality [61]. This film  $WO_{1.92}$  which contains mainly  $W^{4+}$  valence states, exhibits the smallest number of positive ions  $N_{ss}$ , so a small number of oxygen vacancies. This opaque film is close to stoichiometric  $WO_2$  and is not desirable for electrochromic device. Beyond  $Q(O_2) = 1.20$  sccm, the  $\Delta V$  and  $D_{it-average}$  of each  $WO_x$  films increase progressively to reach the highest value when the oxygen flow rate is 3 sccm. This evolution reflects a progressive increase of defects in the volume and a degradation of the interface quality in these films. Moreover, the increase of defects is followed by the formation of a large number of positive ions (increase of  $N_{ss}$ ) due to a large number of oxygen vacancies associated to  $W^{4+}$ ,  $W^{5+}$ , and  $W^{6+}$  valence states. Since these films evolved from semi-transparent to transparent,  $W^{5+}$  are present in the semi-transparent film prepared with  $Q(O_2) = 1.65$  sccm and  $W^{6+}$  are preponderant in the transparent films synthesized with 1.95 sccm and 3 sccm of oxygen flow rate (Table 2). Above  $Q(O_2) = 3$  sccm,  $\Delta V$  which reflects the defects in the volume seems to be less influenced by the increase of  $Q(O_2)$ . On the other hand,  $D_{it-average}$  decreases progressively with the increase of  $Q(O_2)$  showing thus, a progressive reduction of the defects of interfaces with improvement of the quality of the interfaces. The  $N_{ss}$  also gradually decreases, reflecting the decrease of oxygen vacancies in these films. This decrease of oxygen vacancies may explain the decay of  $D_{it-average}$ . The negative value of  $N_{ss}$  for the film prepared with  $Q(O_2) = 14$  sccm is representative of a stoichiometric or over-stoichiometric  $WO_3$ . The  $W^{6+}$  valence states associated with a decrease of oxygen vacancies, is preponderant in the films elaborated with 4.5; 6; 8 and 14 sccm of oxygen flow rate. These transparent films seem to be interesting for electrochromic devices.

Finally,  $Q(O_2) = 3$  sccm seems to be the critical value beyond which we can make transparent  $WO_x$  with nearly constant volume defects whatever  $Q(O_2)$  ( $1.7 < \Delta V < 2$ ). On the other hand, for  $Q(O_2) > 3$  sccm the interface defects ( $D_{it}$ ) and the oxygen vacancies ( $N_{ss}$ ) decrease as the oxygen flow rate increases, improving thus, the quality of the interfaces and the stoichiometry (Table 1). The possibility to vary both the defects and the stoichiometry seems promising in the optimization of the electrochromic properties.

### **3.5 Electrochromic study**

#### **3.5.1 Cyclic voltammetry**

Electrochromic measurements were not possible in our electrochemical windows [-1V; +1V] for films deposited with  $Q(O_2) \leq 3.0$  sccm. This is due to their insufficient content in oxygen, their relatively dense microstructure [62,63] and their opacity. Figures 9 a-d show the first and fifth cyclic voltammograms recorded at 20 mV/s for films deposited with 4.5; 6; 8 and 14 sccm of oxygen flow rate, respectively. The scan was started at 0.25 V (vs. (Saturated Calomel Electrode) SCE) and continued towards negative voltages. The current density is close to zero until the potential drops below - 0.15 V (vs. SCE) from which the current density becomes progressively negative. This cathodic branch corresponds to intercalation, where electrons from the  $SnO_2$  electrode and  $Li^+$  ions from the electrolyte are co-inserted into the  $WO_x$  films. As soon as the current density changes from negative to positive values, after reversing the scan direction from -1.0 V (vs. SCE), deintercalation starts during this anodic branch. The deintercalation current passes through a maximum and finally becomes approximately zero, after reversing the scan direction from +1.0 V (vs. SCE), to a potential return close to 0.25 V (vs. SCE).

During the 1<sup>st</sup> J-V measurement, Figures 9a-d show an important insertion of charge carriers into WO<sub>x</sub> films reflected by an important cathodic current density whose slope of -2 mA/V is nearly constant for all films. But surprisingly, no coloration and discoloration were observed with eyes despite this constant speed of cathodic insertion. For some authors [64,65], the observed phenomenon during the 1<sup>st</sup> J-V measurement for over-stoichiometric films (3.3 < O/W ratio < 3.5) can be attributed to the formation of neutral compound Li<sub>2</sub>O<sub>i</sub> called Lithia. In these films, the excess oxygen is incorporated as interstitials in the structure of the film, and can serve as trapping centers for lithium during the co-insertion of lithium ions and electrons. As far as we are concerned, the films being under-stoichiometric (Table 1), we can rather partially explain this phenomenon by the presence of defects in WO<sub>x</sub> films. These defects are electrically active and can behave like trapping or recombination sites. Therefore, during the first insertion, we think that most of the charge carriers are trapped in these defects so that reduction of W<sup>6+</sup> to W<sup>5+</sup> does not occur. Moreover, as the volume defects in these films are almost constant with approximately the same ΔV (Table 3), then the cathodic insertion rate and the cathodic current density during the first J-V cycle are close for all the films. After several J-V measurements, most of the trapping sites or defects are saturated and the J-V voltammograms become stable from the 5<sup>th</sup> J-V measurement with an electrical reversibility (ER) above 90 %. Therefore, during the cathodic scan, a reduction from W<sup>6+</sup> to W<sup>5+</sup> occurs in transparent WO<sub>x</sub> thin films which become dark blue. Otherwise, during the anodic scan, an oxidation from W<sup>5+</sup> to W<sup>6+</sup> resulting in the original transparent state of the films occurs. Electrochemical reaction allows the coloration/discoloration when both Li<sup>+</sup> and e<sup>-</sup> are electrochemically intercalated/deintercalated in WO<sub>x</sub> thin films

As we can see on the Figures 9 a-d, surfaces delimited by the 5<sup>th</sup> J-V curves are not the same, especially in anodic branch and depend on Q(O<sub>2</sub>). This reflects different number of

charges intercalated and deintercalated for each film. Films indicate a cathodic branch where the current density is negative between - 0.15 V and - 1 V during coloration or reduction process. During discoloration or oxidation process, the anodic branches with positive current density for films deposited with  $Q(O_2) = 4.5$  sccm (Figure 10a) and 6 sccm (Figure 10b) are flat and very spread out up to +1 V (vs. SCE). This reveals a relatively slow discoloration kinetics of these films which seem resisting to  $Li^+$  and  $e^-$  disinsertion. This slow discoloration kinetics and this difficulty in bringing out the charge carriers can be explained by the combined effects of the presence of numerous interface defects revealed by the highest values of the  $D_{it-average}$  (Table 3), by oxygen deficiency reflected by a higher positive ion density ( $N_{ss}$ ) (Table 3) and by their relatively dense microstructure (lower porosity Figure 3).

In comparison, the anodic branches of films elaborated at 8 sccm (Figure 10c) and 14 sccm (Figure 10d) of oxygen flow rate, are rounded and stop quickly before + 0.25 V (vs. SCE). This reveals a relatively rapid discoloration kinetics of these films which seem sensitive to  $Li^+$  and  $e^-$  disinsertion with a rapid discoloration. This sensitivity can be attributed to the combined effects of the decrease of the  $D_{it-average}$ , of oxygen deficiency reflected by  $N_{ss}$  (Table 3) and to their porous microstructure (Figure 3).

### 3.5.2 Chronoamperometry

Figures 10 a-b present the evolution of the current density vs time during insertion (coloration) or disinsertion (discoloration) of charges. All the films were colored during 30 s (Figure 11 a). At the beginning of the coloration process ( $t = 0$  s - 2.5 s), the films exhibit different cathodic current density which increase with oxygen flow rate and revealing that the films close to stoichiometry, the most porous with the least defects exhibit the highest insertion kinetics. This is in accordance with the evolution of porosity, of stoichiometry and defects in

the films as a function of  $Q(O_2)$ . The current density tends to stabilize around  $-0.5 \text{ mA/cm}^2$  towards  $t = 30 \text{ s}$  for all the films, reflecting thus, a slowing down of the co-insertion ( $\text{Li}^+$  and  $e^-$ ) and of the coloring rate. During discoloration, Figure 11b indicates a faster decay in the current's density for films deposited with 8 sccm and 14 sccm of oxygen flow rate compared to films deposited with  $Q(O_2) = 4.5 \text{ sccm}$  and  $6 \text{ sccm}$ . This reveals that the disinsertion rate is favored by a high porosity, a minimum of interface defects and a low deviation from the stoichiometry of the films.

These results suggest that bleaching and coloration kinetics are faster for higher  $Q(O_2)$  films as already observed during cyclic voltammetry measurements.

### 3.5.2.1 Optical modulation

The optical transmittance versus wavelength of  $\text{WO}_x$  thin films at colored and bleached state during the chronoamperometry measurements are shown in Figures 11 a-d for all films deposited with  $Q(O_2) = 4.5, 6, 8$  et  $14 \text{ sccm}$ , respectively. We performed ex-situ optical measurements. As seen, transmission decreases sharply as  $\text{WO}_x$  films are darkened. According to Figure 12, all deposited films are optically reversible. We must precise that, in Figures 11 a-d, dash lines for bleached state represent the initial optical transmittance of as deposited a- $\text{WO}_x$  films. In the Figure 12, we can notice a shift and superposition of the absorption threshold of colored and bleached films toward smaller wavelengths reflecting a modification of their bandgap after intercalation/deintercalation of charges. Optical modulation ( $\Delta T$ ), which is the difference of optical transmittance between the bleached and the colored states of films is also shown on this Figure 12 for  $\lambda = 632 \text{ nm}$  [66,67]. The  $\Delta T$  of films are very close after 30 s of insertion. It can be explained by the fact that the current density reached after 30 s during coloration process is close to  $0.5 \text{ mA/cm}^2$  for all the films, despite different kinetics of insertion. On the other hand, to reach the optical reversibility during the disinsertion of the films ( $\Delta T$  discoloration), the times required are very different, revealing slower discoloration kinetics for the lowest  $Q(O_2)$ .



### 3.5.2.2 Coloration efficiency (CE)

One quantitative way to characterize the color change in electrochromic materials is to calculate their coloration efficiency (CE). The CE is then, an important index to follow during the fabrication of electrochromic devices. This parameter is given by [68,69]:

$$CE = \frac{\Delta OD}{Q_{in}} \quad (\text{eq.8})$$

Where CE ( $\text{cm}^2/\text{C}$ ) is the coloration efficiency,  $Q_{in}$  ( $\text{C}/\text{cm}^2$ ) is the injected charges quantity into the film per unit area and  $\Delta OD$  is the optical density change. The optical density means the light absorption ability of electrochromic films during their 2 extremes optical states.  $\Delta OD$  is calculated with the following equation:

$$\Delta OD = \log\left(\frac{T_b}{T_c}\right) \quad (\text{eq.9})$$

where  $T_b$  and  $T_c$  are the transmissions in the bleached and the colored states, respectively. A high CE indicates that the material structure exhibits a large optical modulation with a small charge insertion. The evolution of CE versus wavelength is given in Figure 13. It can be seen from the Figure 13 that the coloration efficiency of  $\text{WO}_x$  films decreases with increasing oxygen flow rate and tend to increase vs wavelength in the visible region, for all the films. The maximum coloration efficiency is reached at 770 nm. The decay of the CE for films prepared with 8 sccm and 14 sccm of oxygen flow rate, could be due to their higher porosity close to 20 %, but mainly to the decrease of oxygen deficiency. Probably, their porous microstructure does not enable availability of enough insertion sites for  $\text{Li}^+$  ions intercalation to form the tungsten bronze. The coloration efficiency of these films at 632 nm are  $33.0 \text{ cm}^2/\text{C}$  and  $34.3 \text{ cm}^2/\text{C}$ ,

respectively. The highest CE values are found to be 44.4 cm<sup>2</sup>/C and 42.2 cm<sup>2</sup>/C at 632 nm for respectively Q (O<sub>2</sub>) = 4.5 sccm and 6 sccm, P<sub>t</sub> = 1.13 Pa and 1.37. These WO<sub>x</sub> films CE is improved firstly because of their less porous microstructure (Figure 3) as it was also observed by Sun et al. [15]. Their less porous microstructure may allow the availability of Li<sup>+</sup> insertion sites in order to generate a large number of W<sup>5+</sup>. Otherwise, the main reason for having highest CE with these films is their large number of positive N<sub>ss</sub> associated to oxygen vacancies (Table 3). According to Lee et al. [18], the CE increases with increasing oxygen vacancies in WO<sub>x</sub> films and this permit to generate a great number of W<sup>5+</sup> by insertion of Li<sup>+</sup> and e<sup>-</sup>. As shown in Figure 14, showing the evolution of the electrical parameters and the CE at 632 nm as a function of Q(O<sub>2</sub>), one can note an interesting correlation between the interface defects (D<sub>it</sub>) and oxygen deficiency (N<sub>ss</sub>) in the films and their coloration efficiency (CE). Other authors confirm that CE changes can be guided by controlled manipulation of oxygen vacancies in films of perovskite structure without additional Li<sup>+</sup> ion involvement [70].

These films, with improved CE, have the highest D<sub>it-average</sub> and N<sub>ss</sub>. As a consequence, we can assume that the control of these electrical parameters can contribute to the enhancement of the CE of WO<sub>x</sub>.

### 3.5.3 Response time

The response time, which measure the time interval necessary for switching between the two extreme optical states of films, were determined by the CA. In our study, the coloration time t<sub>c</sub> correspond in the time required to reach a variation in optical density ΔOD = 0.3 at the wavelength of 632 nm. The ΔOD = 0.3 correspond to a decay of 50 % of the transmission during coloration. The discoloration time t<sub>d</sub> corresponds in the time required to return to the

initial discoloration state.  $t_d$ , is determined when the current density begin to stabilized and is close to zero (Figure 11). These response times are summarized in Table 4.

All the films of this study deposited with  $Q(O_2) \geq 4.5$  sccm show staining times substantially close to 20-30 seconds to reach  $\Delta DO = 0.3$ . On the other hand, the discoloration times are very different from 180 to 8 seconds and seem to be strongly dependent both on interface defects and on the density of positive ions (or oxygen deficiency) present in the  $WO_x$  films. Films deposited with 8 sccm and 14 sccm of oxygen flow rate, exhibit short response time de 8 - 9 s contrarily to those deposited with  $Q(O_2) = 4.5$  sccm and 6 sccm, 180 and 75 s respectively. The response time confirm observations made in section 3.4.1 and 3.4.2.

#### 4 Conclusion

This work permitted us to show the effect of oxygen flow rate and total pressure on physical properties of  $WO_x$  thin films. The correlation between these properties and electrochromic performances of  $WO_x$  was clearly illustrated.

The  $WO_x$  optical transmittance evolves progressively from opaque to transparent when films become rich in oxygen. Also, the optical transmittance has been useful to investigate on the porosity of transparent films. Transparent films deposited with a high total pressure were porous with degraded refractive index whereas transparent films deposited with low total pressure where less porous with high refractive index. Thanks to Raman spectroscopy, we have been able to prove the existence of different W valence states and the link between transparency and the  $W^{6+}$  valence state. The  $C_{ox}$  and  $D_{it}$  of the films decrease with increasing  $Q(O_2)$  and  $P_t$ .

Also, the  $N_{ss}$  decreases with increasing  $Q(O_2)$  and tend to become negative with an excess of  $O_2$  in the films.

Electrochromic performances study was possible for films deposited with  $Q(O_2) \geq 4.5$  sccm and  $P_t \geq 1.13$  Pa, in an electrochemical windows [-1 V; +1V]. Results show that the lowest values of  $\Delta T$ , CE and a rapid insertion/disinsertion kinetics were achieved for films deposited with high ( $Q(O_2)$ ;  $P_t$ ). These films are porous with low refractive index. They exhibit a low  $D_{it-average}$  and low positive  $N_{ss}$ , or a negative  $N_{ss}$ .

Otherwise, highest values of  $\Delta T$  and CE, with slow insertion/disinsertion kinetics were achieve for films deposited with low ( $Q(O_2)$ ;  $P_t$ ). These films are slightly less porous with high refractive index. They exhibit higher values of  $D_{it-average}$  and positive  $N_{ss}$ . A reduction of a large number of  $W^{6+}$  to  $W^{5+}$  is supposed to be formed during a double insertion of  $Li^+$  and  $e^-$  in these films.

Finally, we have shown in this study, a correlation between the electrical parameters, the optical properties and the electrochromic performances. Since the  $D_{it-average}$  reflects defects at interface and  $N_{ss}$  the oxygen deficiency, we think that a good control of these parameters, is of a great interest for the manufacturing of electrochromic devices with optimized coloration efficiency.

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## Tables list

Q (O <sub>2</sub> ) (sccm)	Total pressure (Pa) ± 0.01Pa	Films thickness (nm) ± 15 nm	EDS ratio (at. % O /at. % W)
1.20	0.45	589	1.92
1.65	0.53	582	2.56
1.95	0.61	565	-
3.00	0.83	507	2.70
4.50	1.13	500	2.85
6.00	1.37	502	2.92
8.00	1.70	528	2.94
14.00	2.60	487	3.00

Table 1: Thicknesses and EDS ratio of WO<sub>x</sub> thin films deposited onto glass with various oxygen flow rates and total pressures. The relative error on EDS ratio is 3%.

Attribution	Q(O <sub>2</sub> ) (sccm)			Attribution
	14 → 1.95	1.65	1.20	
O-W <sup>6+</sup> -O bending [43]	278 (140)	221 (93)		W <sup>4+</sup> -O [45]
		268 (35)		W <sup>4+</sup> -O
		320 (114)		O-W <sup>5+</sup> -O [45]
		423 (71)		W <sup>5+</sup> =O [45]
		458 (26)		
O-W <sup>6+</sup> -O stretching [39,43]	671 (122)	671 (122)		O-W <sup>6+</sup> -O stretching [39,43]
	782 (153)	782 (153)		
		922 (145)	885 (115)	stretching
W <sup>6+</sup> =O stretching [39]	961 (95)	961 (95)		W <sup>6+</sup> =O stretching [39]
	1015 (20)	1015 (20)		

Table 2 : Wavenumbers and Full Width at Half Maximum (in parenthesis), cm<sup>-1</sup>, of the bands deduced from the refinements (gaussian profile) of the Raman spectra of samples deposited with Q(O<sub>2</sub>) ≥ 1.95 sccm, Q(O<sub>2</sub>) = 1.65 sccm, and Q(O<sub>2</sub>) = 1.20 sccm.

Q (O <sub>2</sub> ) (sccm)	Hysteresis $\Delta V$ (V)  $\pm 2 \%$	D <sub>it-average</sub> * 10 <sup>12</sup> states/cm <sup>2</sup> /eV  $\pm 7.5 \%$	Ions density*10 <sup>12</sup> N <sub>ss</sub> (ions/cm <sup>2</sup> )  $\pm 7.6 \%$	Oxide capacitance * 10 <sup>10</sup> (Farad)  $\pm 1.2 \%$
1.20	0.003	1.1	0.9	8,88
1.65	1.43	4.5	11.2	6,06
1.95	1.80	16.4	14.2	5,75
3.00	2.80	35.7	12.4	5,52
4.50	2.00	24.4	10.7	5,45
6.00	1.72	9.2	4.1	6,07
8.00	2.00	3.9	1.1	5,36
14.00	1.70	5.3	-1.8	4,01

Table 3: Oxygen flow rates with the various corresponding defects identified in the film by the measurements of the electrical parameters  $\Delta V$ , D<sub>it</sub>, N<sub>ss</sub> and C<sub>ox</sub>. The relative errors on the hysteresis mid-height width, interface traps average distribution, ions density and oxide capacitance are 2 %, 7.5 %, 7.6 % and 1.2 %, respectively.

Q O <sub>2</sub> (sccm)	t <sub>coloration</sub> (s)	t <sub>discoloration</sub> (s)
4.5	19	180
6	22	75
8	30	9
14	27	8

Table 4: Response time for WO<sub>x</sub> films.



## Figures caption

Figure 1: (a) Effect of oxygen flow rate on optical transmittance spectra as a function of the wavelength of  $\text{WO}_x$  films deposited on glass substrate. (b) Evolution of average optical transmittance in the visible region as a function of oxygen flow rate. In insert, XRD patterns of a- $\text{WO}_x$  films deposited on glass for  $Q(\text{O}_2) = 1.2, 4.5$  and  $14$  sccm.

Figure 2: Influence of oxygen flow rate increase and total pressure increase on the evolution of  $\text{WO}_x$  thin films refractive index as a function of the wavelength in visible region.

Figure 3: Evolution of average refractive index between  $425 \text{ nm} \leq \lambda \leq 800 \text{ nm}$ . and average porosity of transparent  $\text{WO}_x$  thin films as a function total pressure and oxygen flow rate.

Figure 4: SEM observations were performed on cross sections of the films deposited on silicon substrates with oxygen flow rates of  $1.2$  sccm ( $\text{WO}_{1.92}$ ),  $4.5$  sccm ( $\text{WO}_{2.85}$ ) and  $14$  sccm ( $\text{WO}_3$ ).

Figure 5: Raman spectra of selected samples deposited with  $Q(\text{O}_2) = 14.00, 1.95$  and  $1.20$  sccm (accumulation time  $1200$  s,  $10$  accumulations), and Raman spectrum of crystallized  $\text{WO}_3$  (monoclinic) collected in the same spectral range (accumulation time  $120$  s,  $2$  accumulations). The single stars (\*) stands for contamination by very narrow plasma lines. The double stars (\*\*) indicate contamination by the set of narrow Raman lines of air (rotational lines) in the near vicinity of the surface of the samples. The dotted lines define the low and high frequency ranges selected for the refinements, as already considered [39,43].

Figure 6: Raman spectra of samples deposited with  $Q(\text{O}_2) = 1.95, 1.65,$  and  $1.20$  sccm with their fitted curves (colored lines) obtained by least squares method. Each spectrum is

deconvoluted by the sum of gaussian shape bands. Stars (\*) at  $235\text{ cm}^{-1}$  stands for contamination by plasma line.

Figure 7: Integrate intensity ratio as a function of  $Q(\text{O}_2)$  (a), and as a function of the EDS O/W ratio (b), with respect to the O- $\text{W}^{6+}$ -O stretching band at  $671\text{-}782\text{ cm}^{-1}$ , of the band at  $278\text{ cm}^{-1}$  attributed to O- $\text{W}^{6+}$ -O bending (square), and the band at  $961\text{ cm}^{-1}$  attributed to  $\text{W}^{6+}=\text{O}$  stretching (dot). The dotted lines are just guides for the eye.

Figure 8: Integrate intensity ratio as a function of  $Q(\text{O}_2)$  (a), and as a function of the EDS O/W ratio (b), with respect to the O- $\text{W}^{6+}$ -O band at  $671\text{-}782\text{ cm}^{-1}$ , of the lines attributed to  $\text{W}^{6+}$  at  $278\text{ cm}^{-1}$  (square), O- $\text{W}^{5+}$ -O at  $320\text{ cm}^{-1}$  (blue dot),  $\text{W}^{5+}=\text{O}$  at  $423$ , and  $458\text{ cm}^{-1}$  (green dot),  $\text{W}^{4+}$  at  $221\text{ cm}^{-1}$  (triangle), and of the narrow line at  $268\text{ cm}^{-1}$  (red diamond). The dotted lines are just guide for the eye.

Figure 9: Hysteresis loop of the  $\text{WO}_x$  film deposited with  $Q(\text{O}_2) = 8\text{ sccm}$

Figure 10: Cyclic voltammograms of  $\text{WO}_x$  films deposited with various oxygen flow rates. The 1<sup>st</sup> and the 5<sup>th</sup> CV curves are shown for each film. The scanning speed is  $20\text{ mV / s}$ .

Figure 11: Evolution of the current density vs time during coloration (a) and discoloration (b).

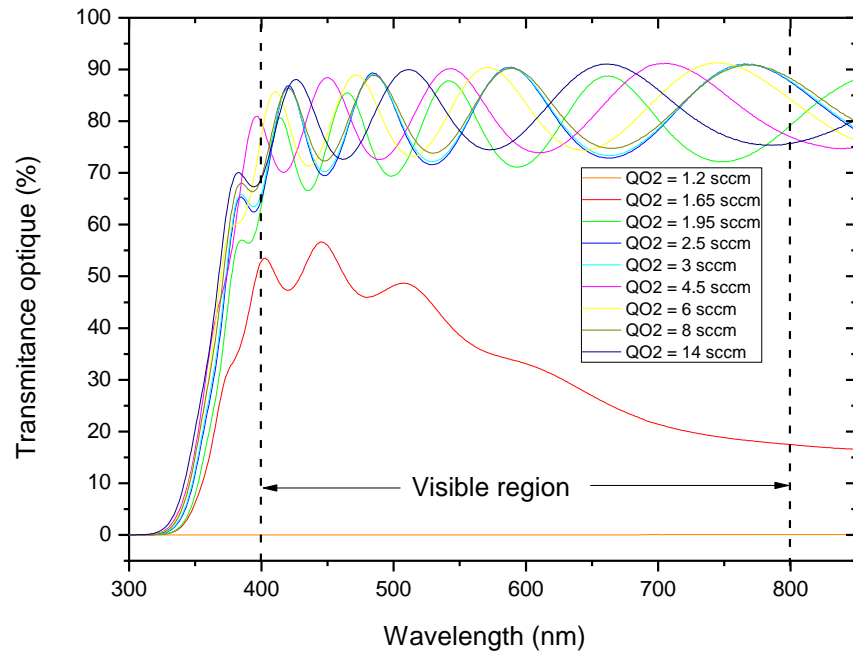
Figure 12: Optical transmittance of as deposited, bleached and colored  $\text{WO}_x$  thin films and optical modulation at  $632\text{ nm}$ . The dash lines for bleached states represent the initial optical transmittance of the as deposited a- $\text{WO}_x$  films.

Figure 13: Coloration efficiency versus wavelength for films deposited with variable oxygen flow rate.

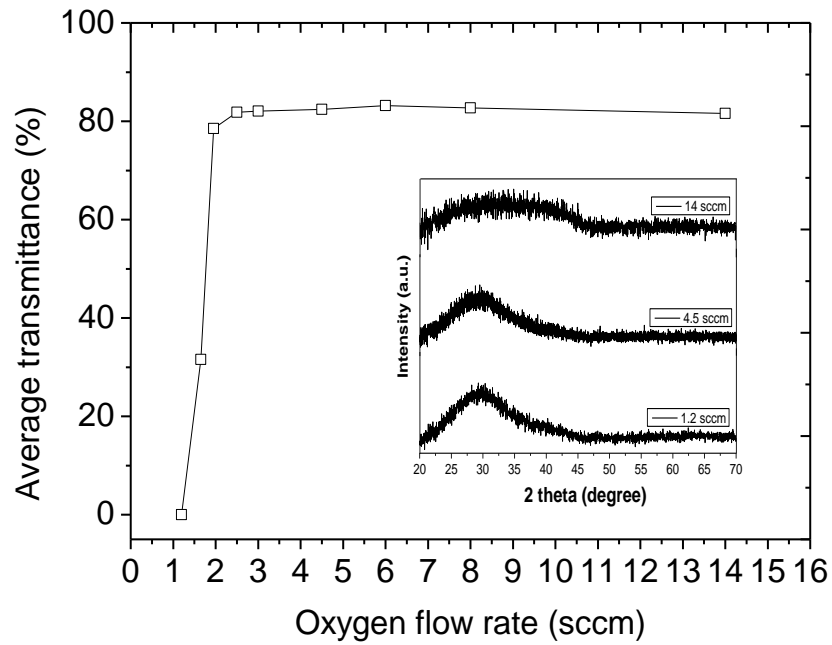
Figure 14: Evolution of electrical parameters and coloration efficiency at 632 nm as a function of oxygen flow rate.



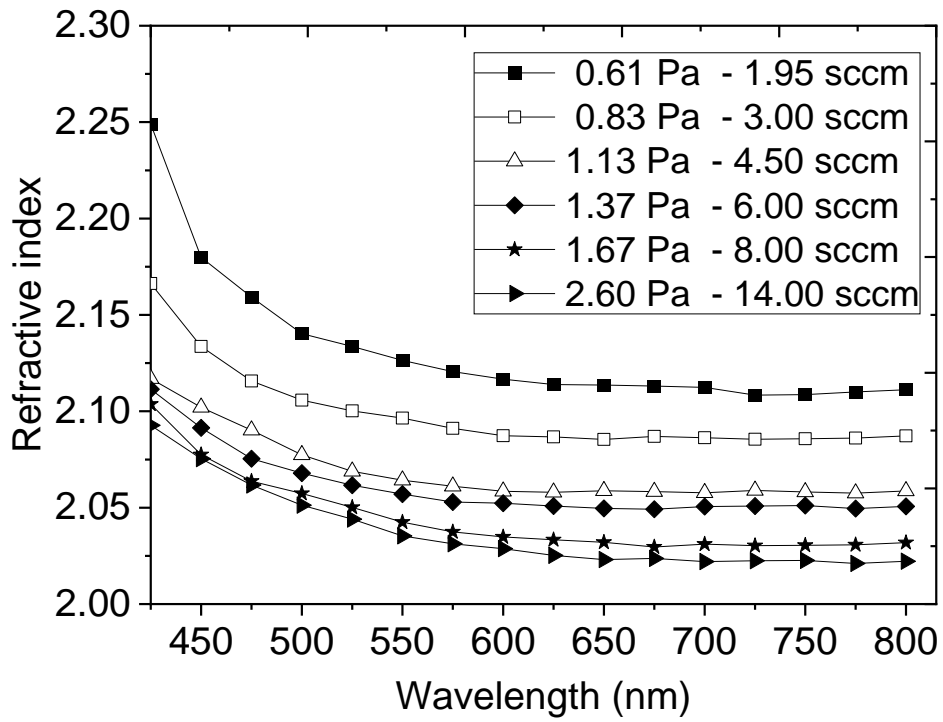
## List of Figures



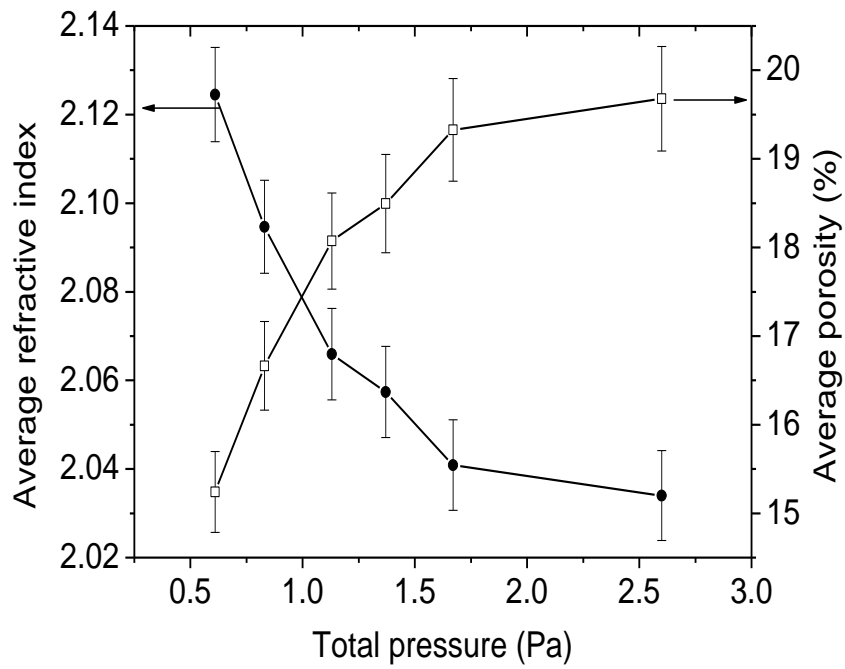
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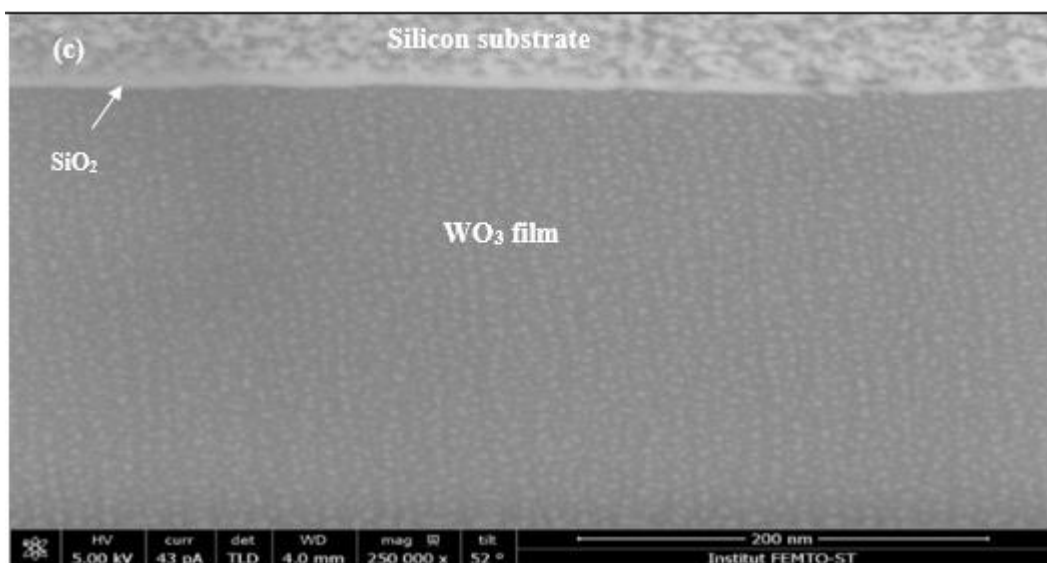
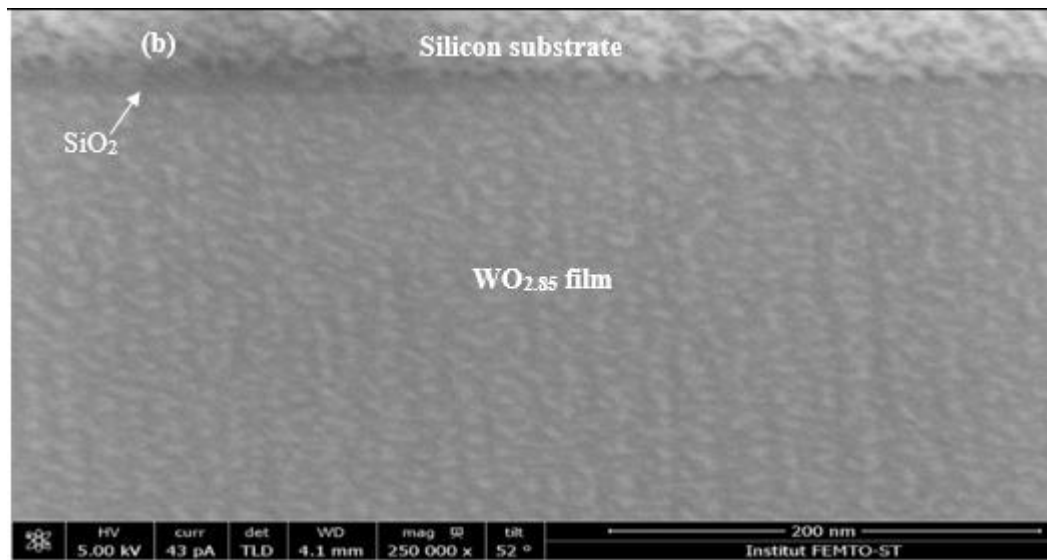
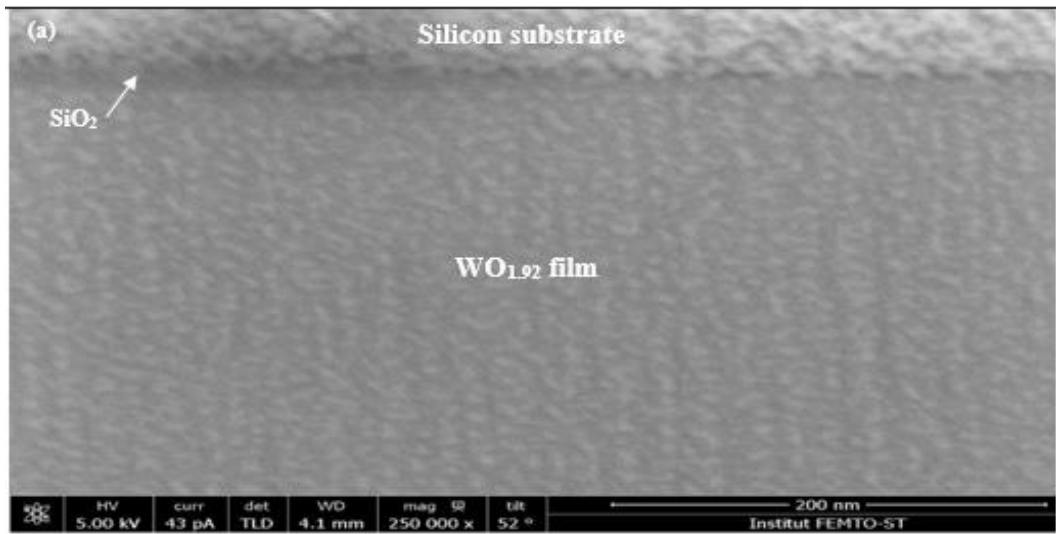
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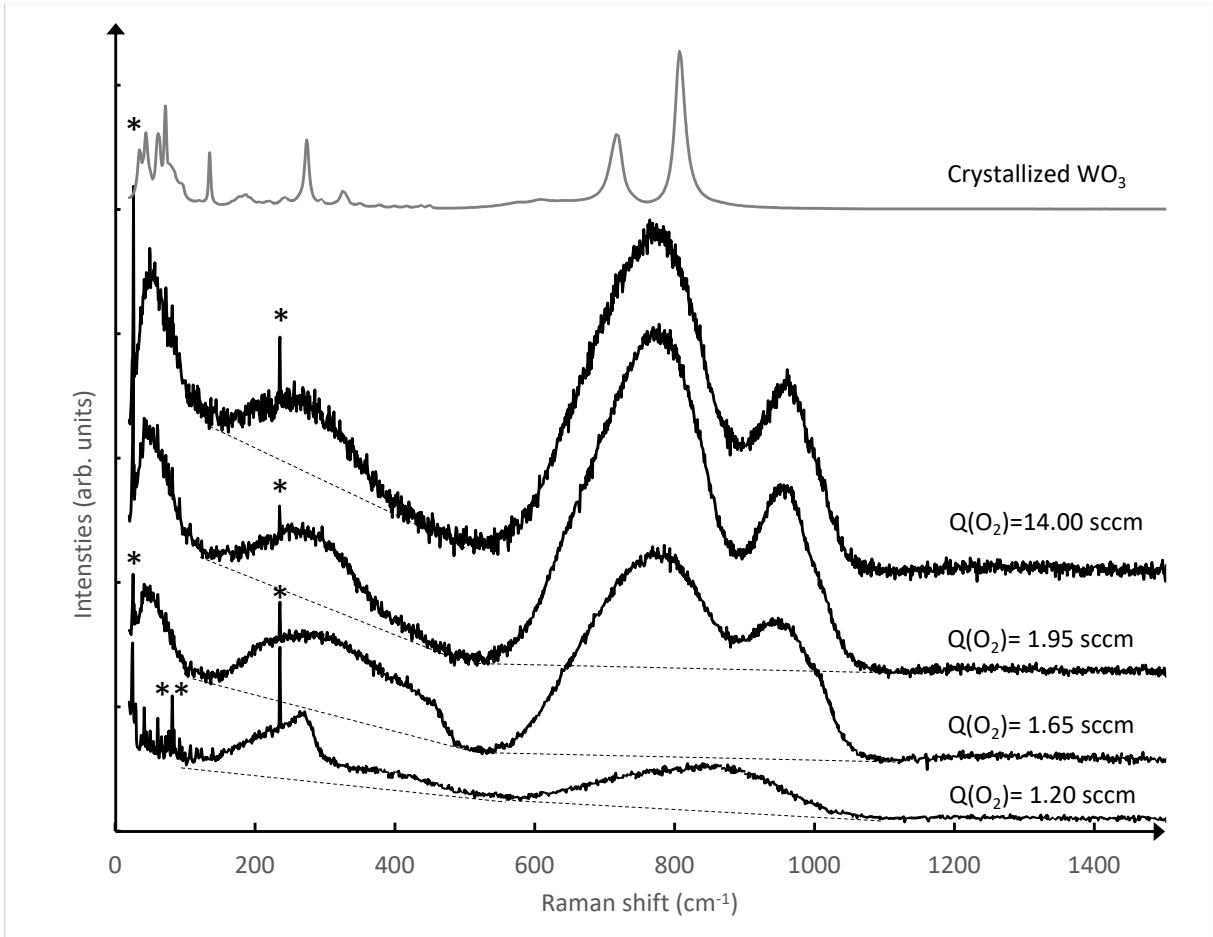


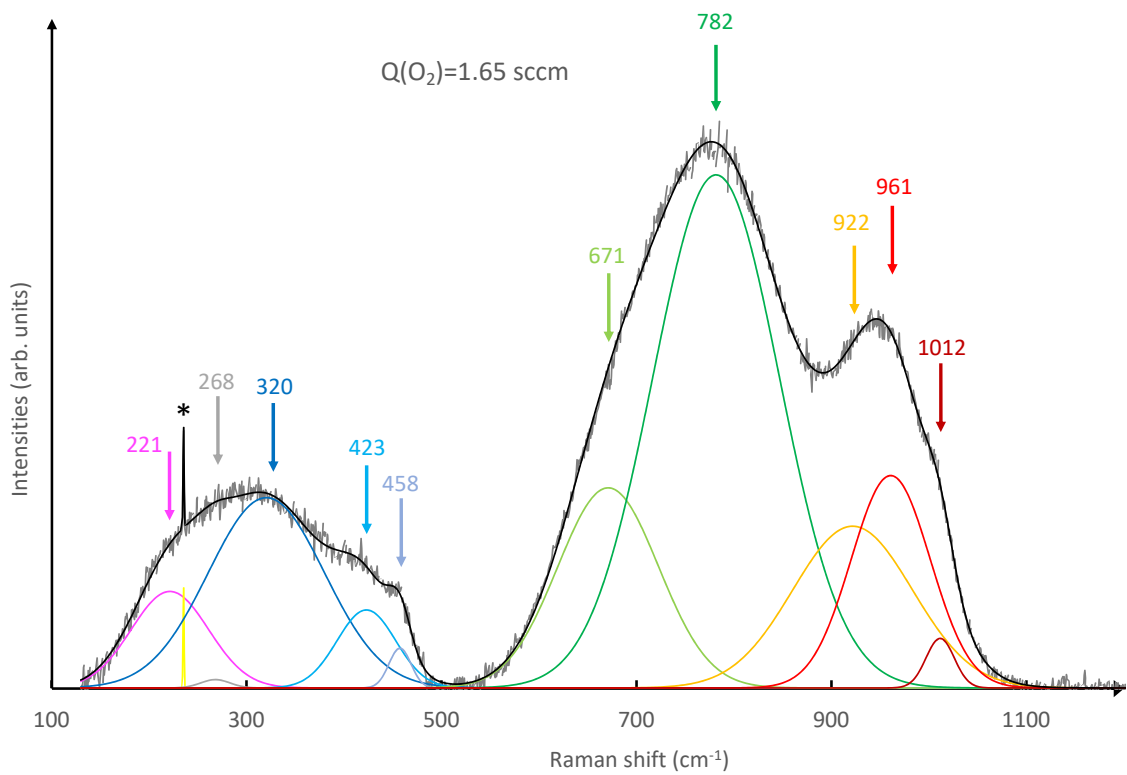
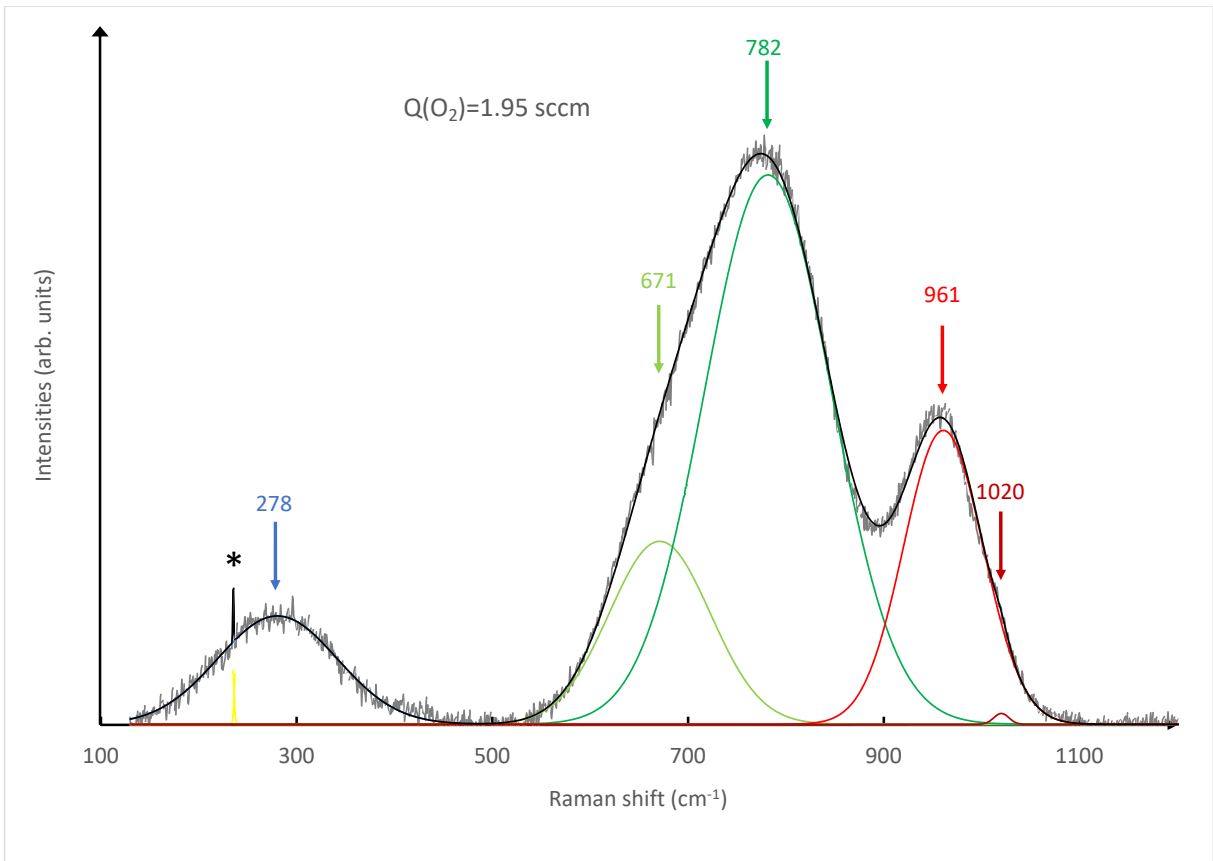
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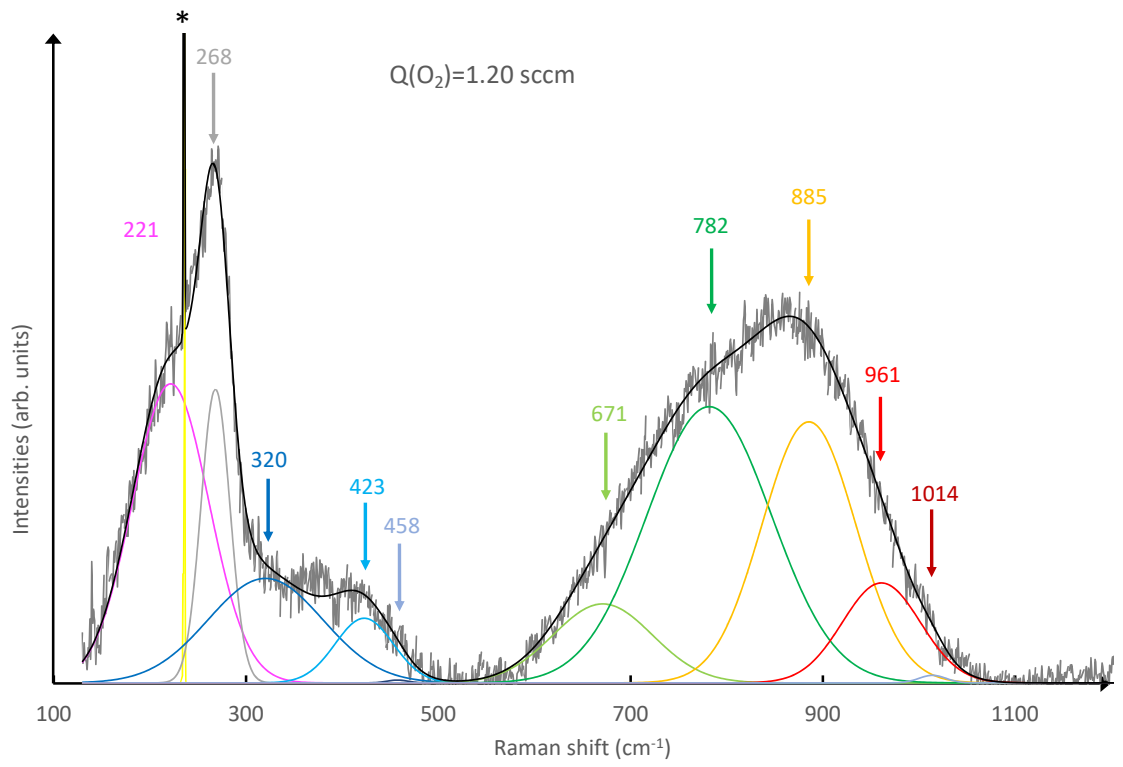


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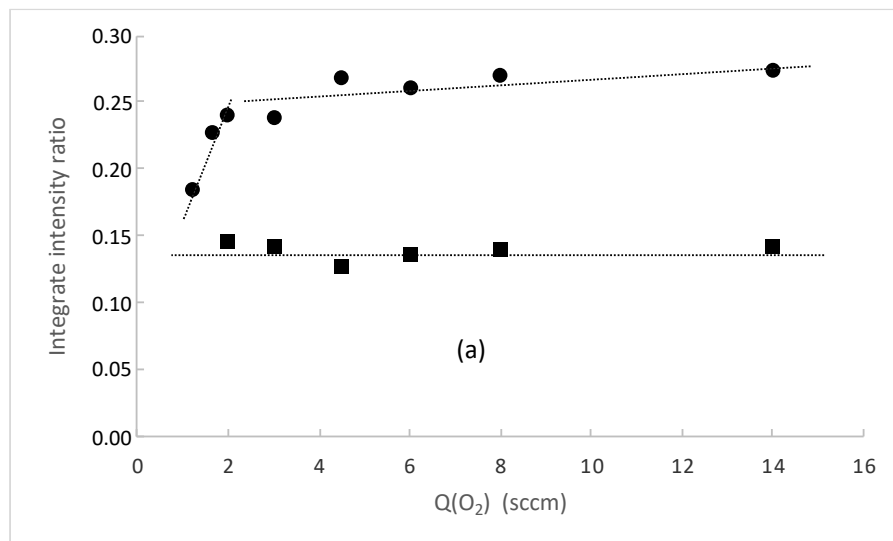




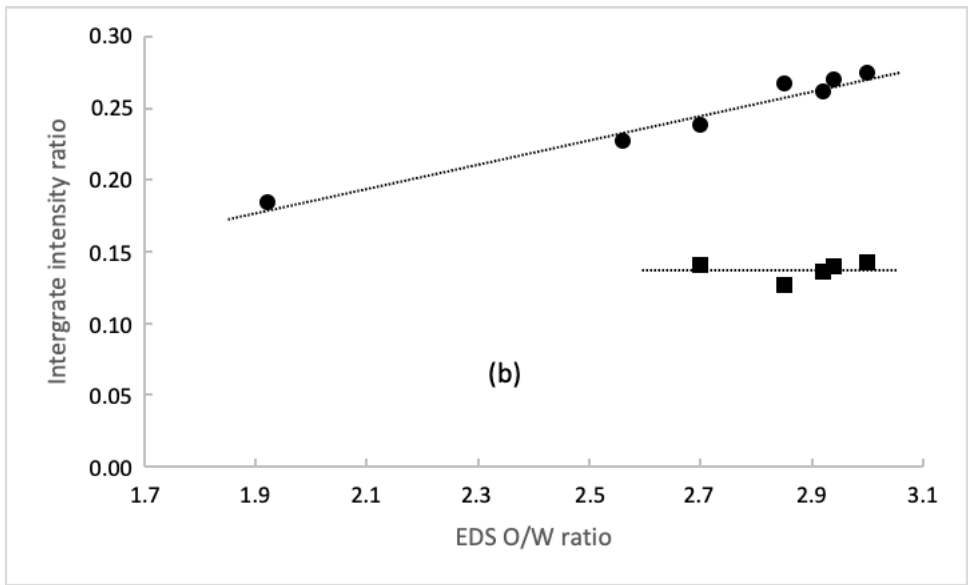




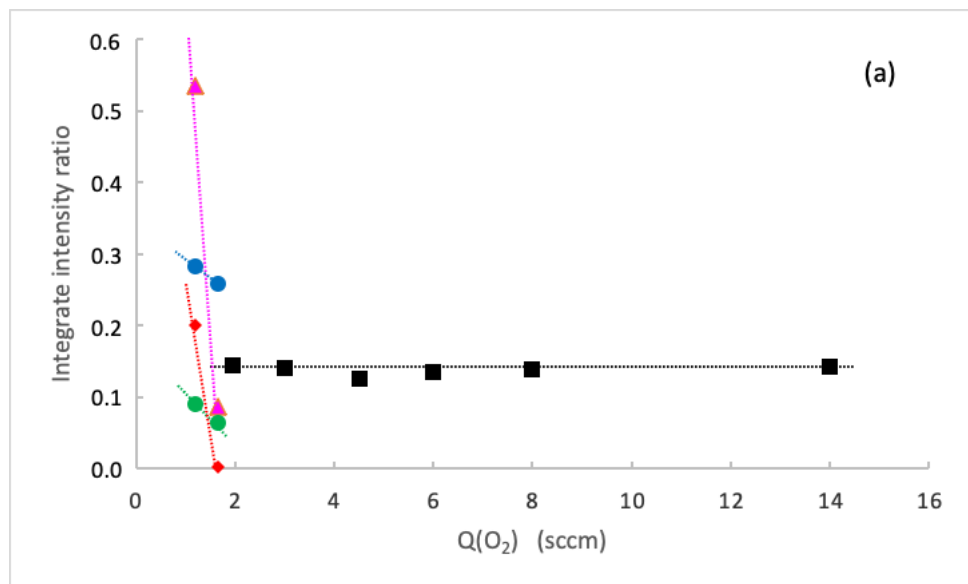
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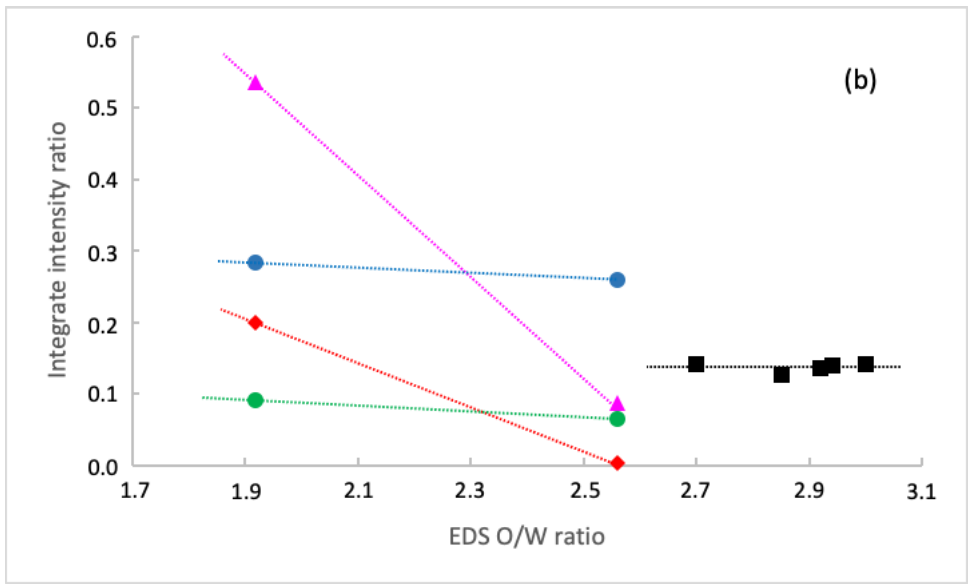
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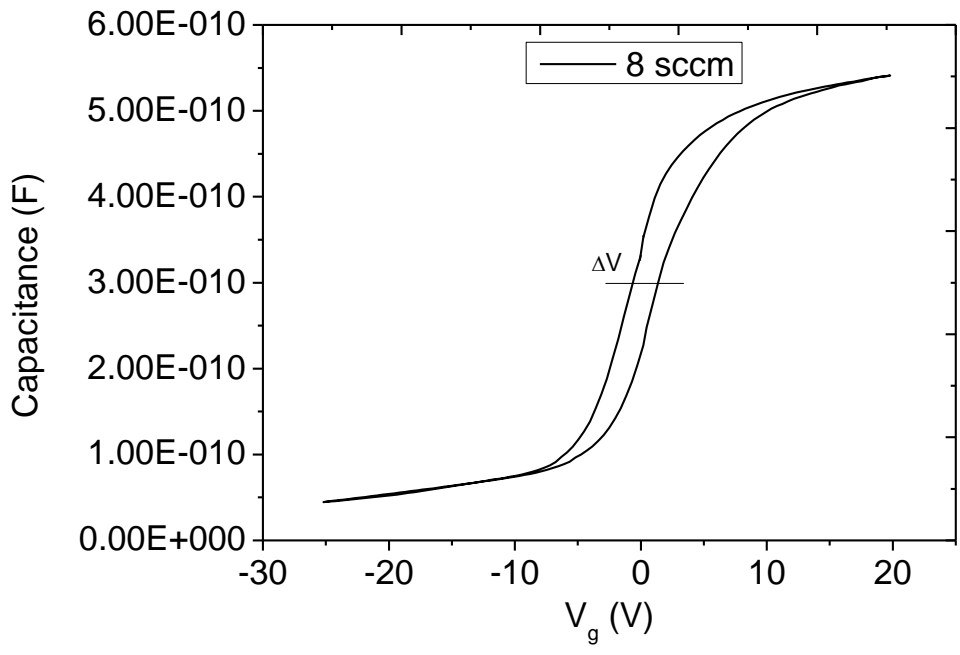
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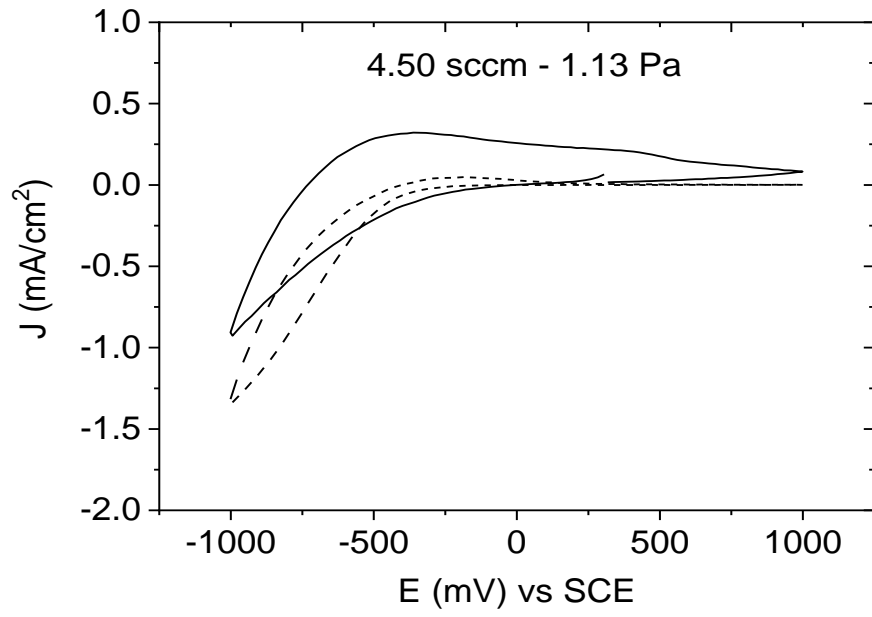


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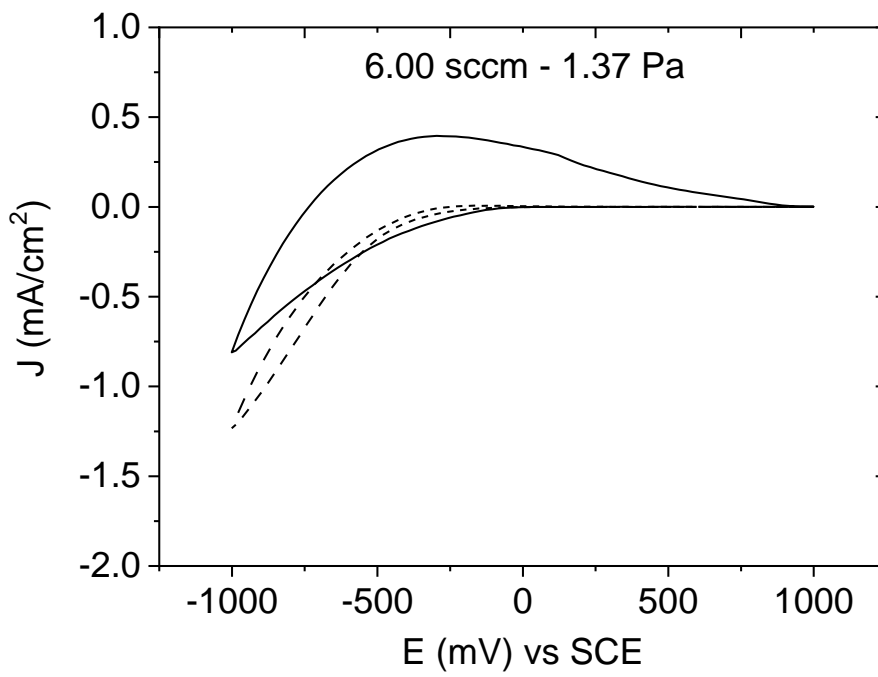


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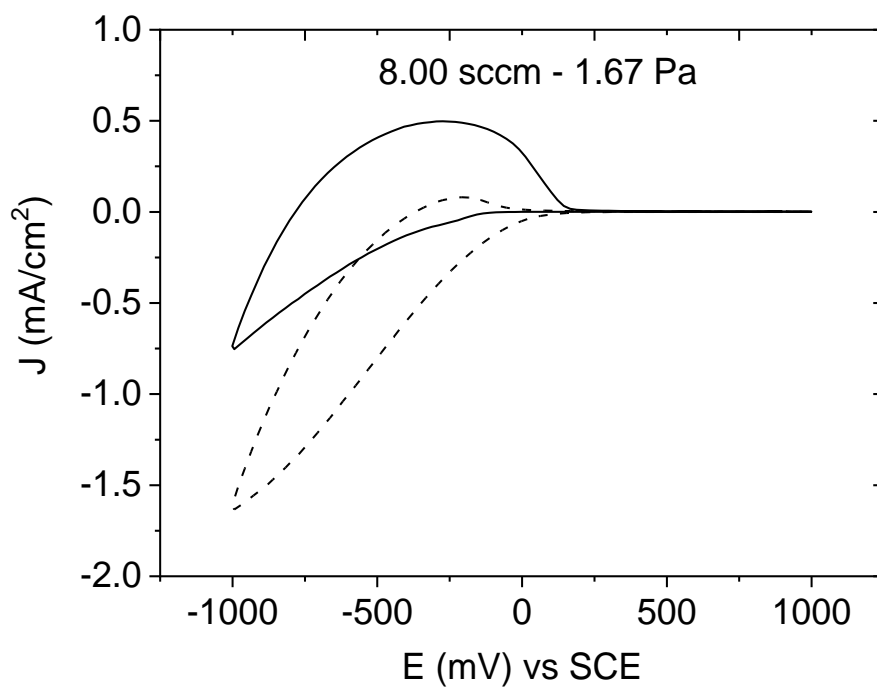




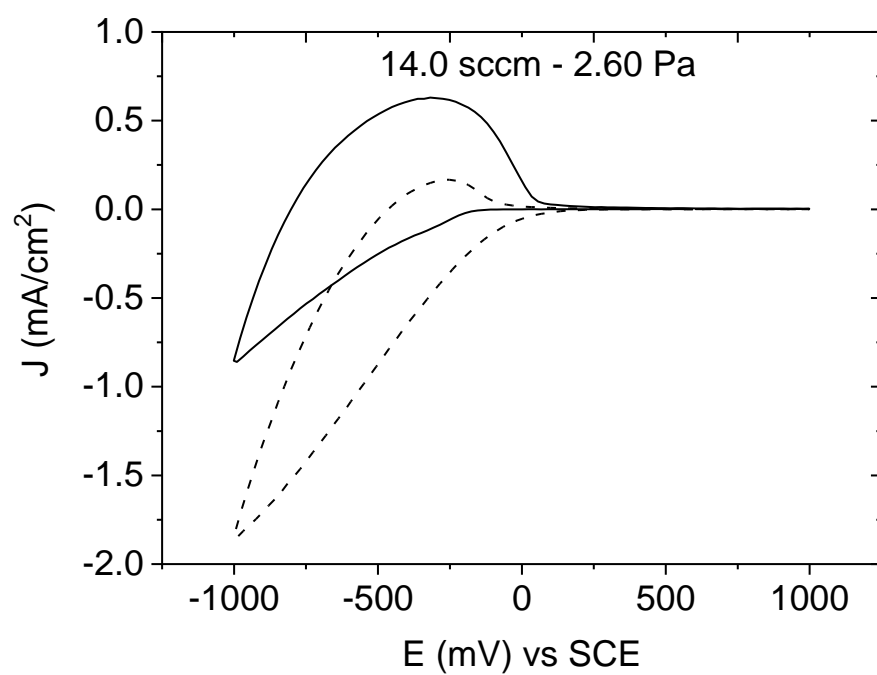
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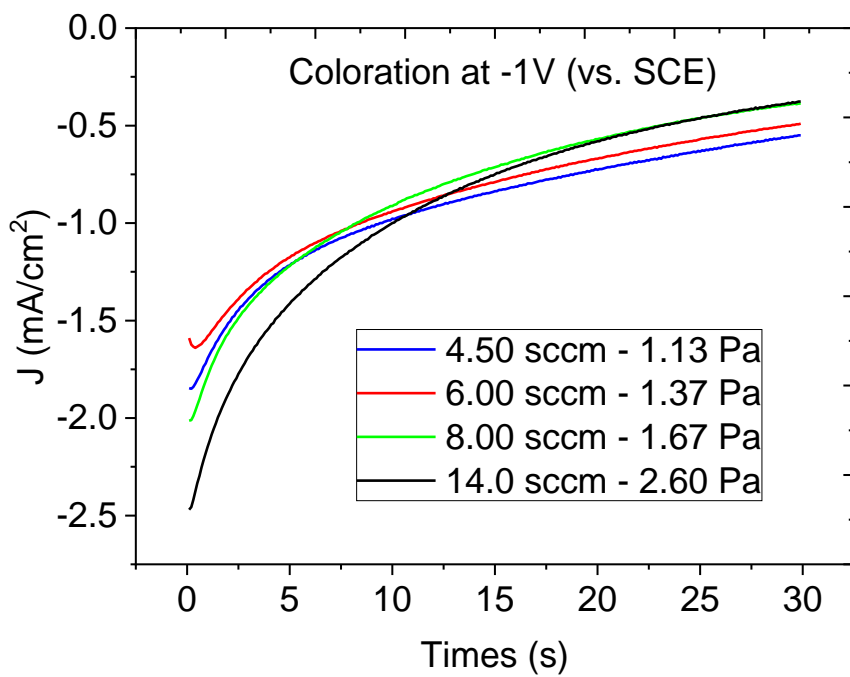
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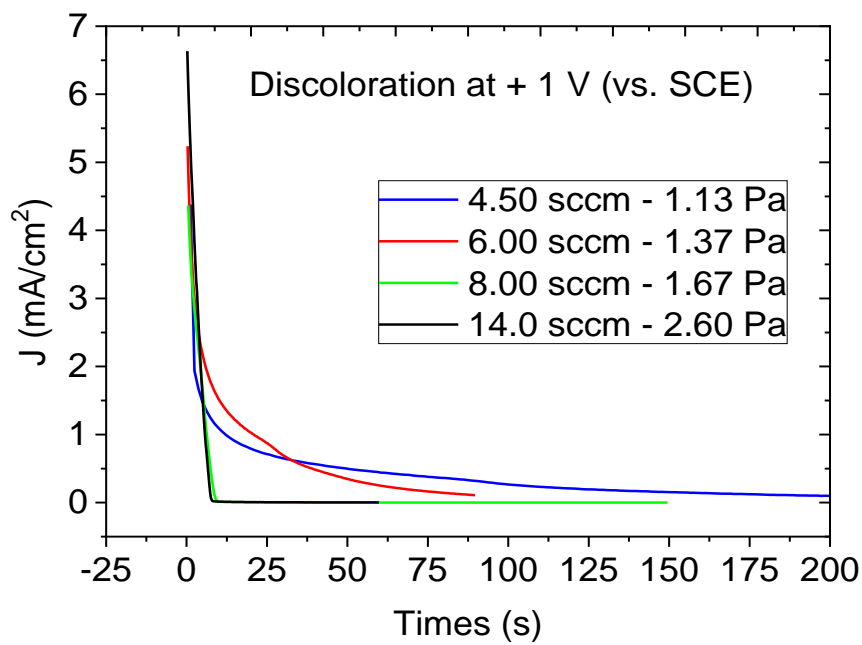
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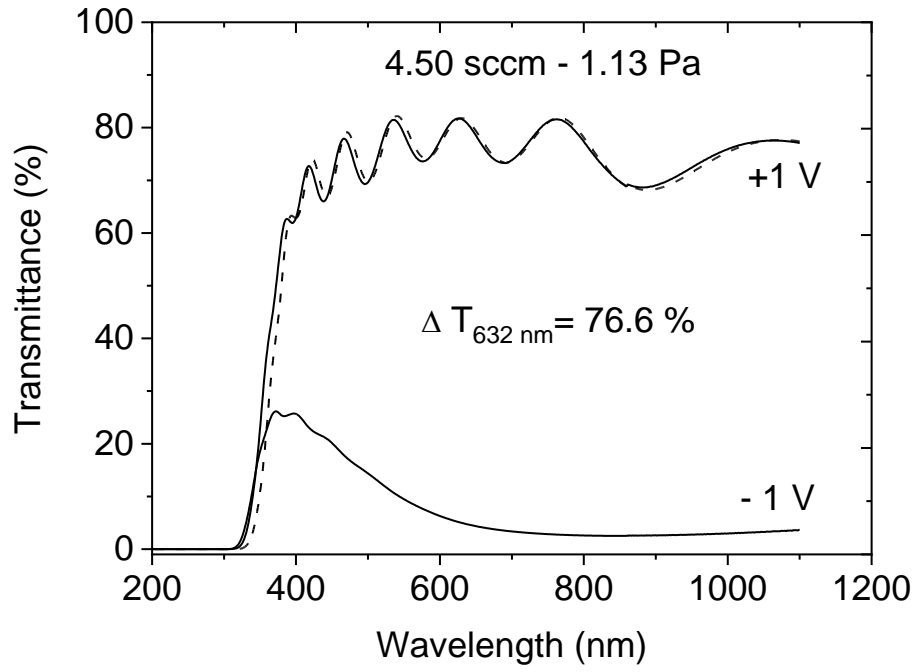
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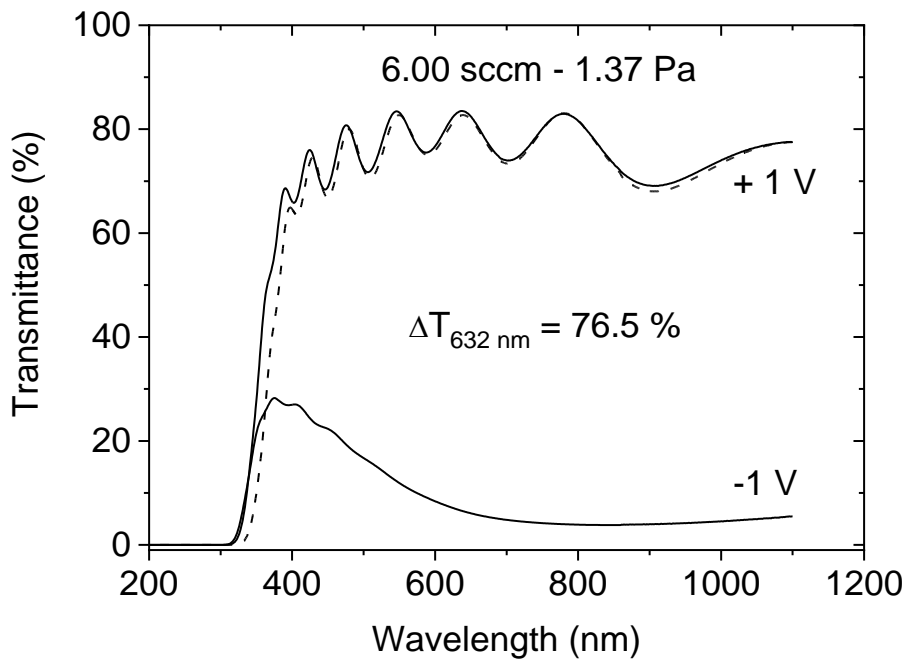
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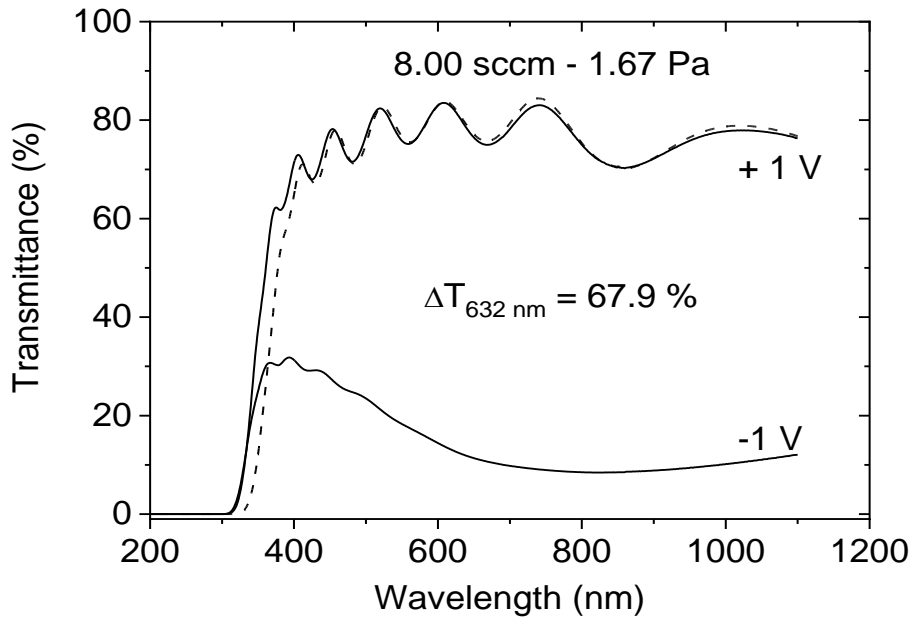
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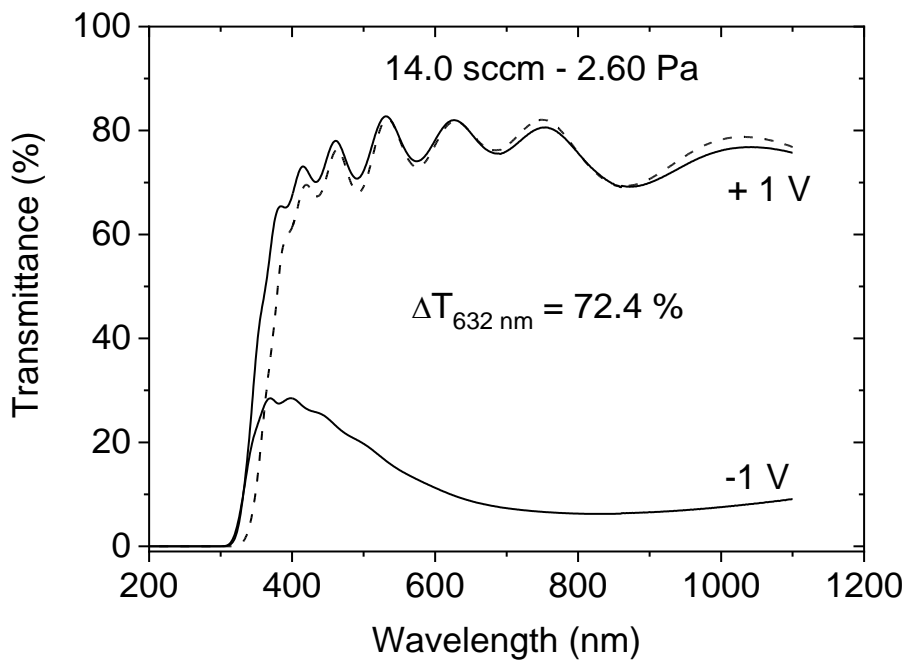
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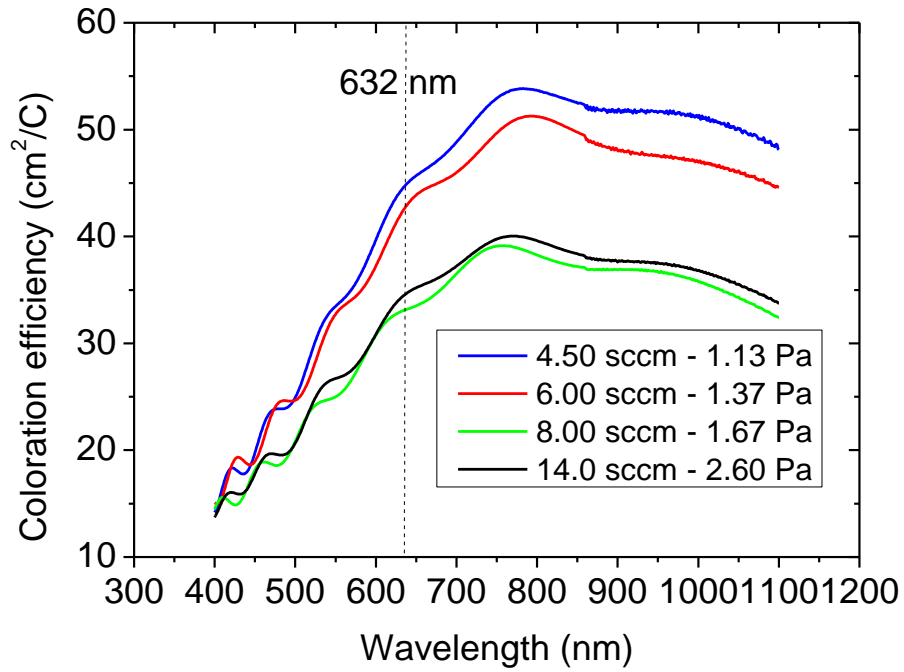
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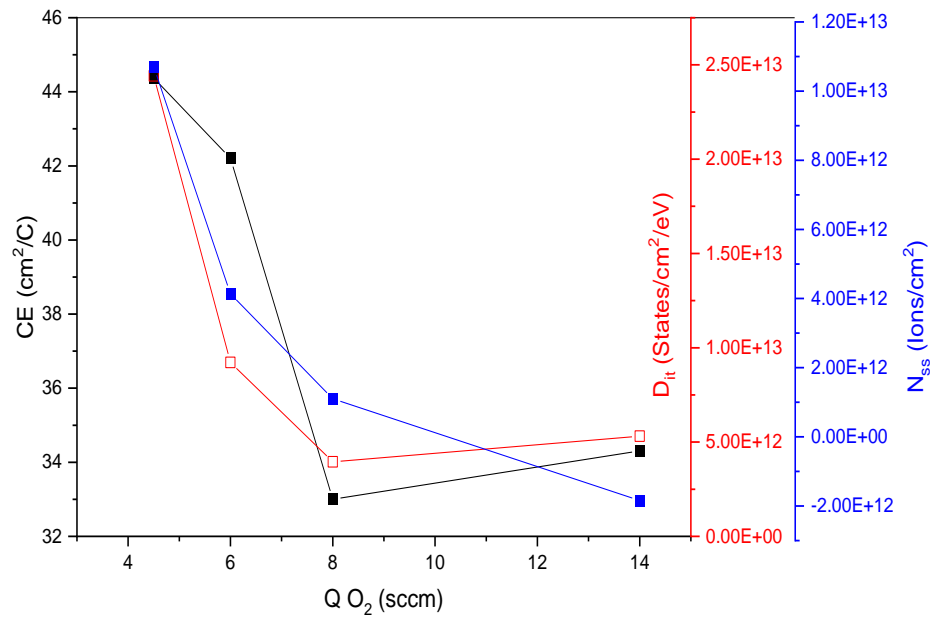
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12 d



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14